

PWA of Baryon Resonances at BES

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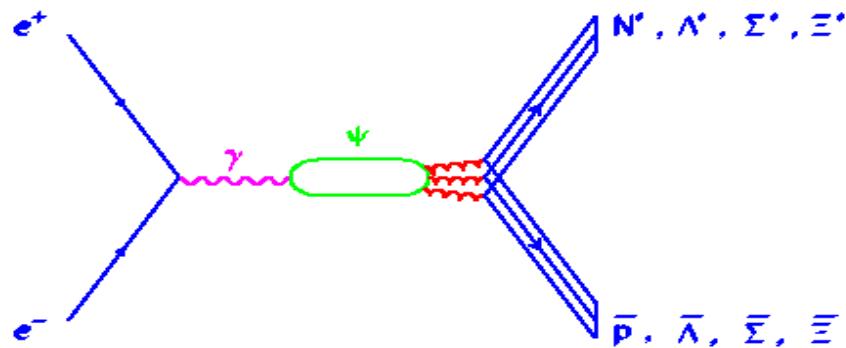
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Outline:

- **Introduction**
- **Event-based PWA procedure**
- **N^* resonances observed in $\psi \rightarrow \bar{N} \pi N, \bar{N} \eta N$**
- **N^* and Λ^* observed in $J/\psi \rightarrow \bar{\Lambda} K N$**
- **Summary and Prospects**

1. Introduction

$$J/\Psi \rightarrow \bar{B}BM \Rightarrow N^*, \Lambda^*, \Sigma^*, \Xi^*$$

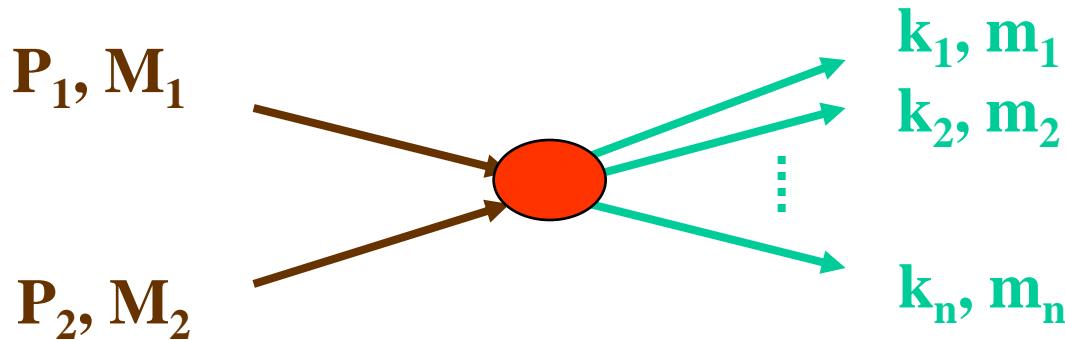


New mechanism for baryon production & an ideal isospin filter

high statistics extension to ψ' , χ_{cJ} , η_c

2. Event-based PWA framework

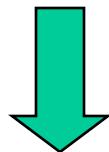
1) Experimental observables – Starting point for PWA



Experimental distribution probability:

$$W_{\text{exp}}(k_1, \dots, k_n) \sim |\mathcal{M}(k_1, \dots, k_n)|^2 \varepsilon(k_1, \dots, k_n) d\Phi_n(k_1, \dots, k_n)$$

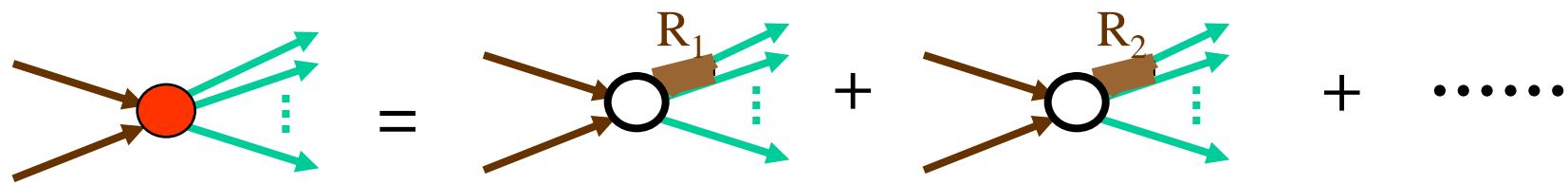
amplitude efficiency phase space



$$d\Phi_n(P; p_1, \dots, p_n) = \delta^4 (P - \sum_{i=1}^n p_i) \prod_{i=1}^n \frac{d^3 p_i}{(2\pi)^3 2E_i}$$

various projections

2) Event-based multi-dimensional PWA fit



$$\mathcal{M}(k_1, \dots, k_n) = C_1 A_{R1}(k_1, \dots, k_n) + C_2 A_{R2}(k_1, \dots, k_n) + \dots$$

→ $W_{th}(k_1, \dots, k_n; C_1, C_2, \dots)$

PWA: fitting experimental data to get C_1, C_2, \dots

Old fashion : χ^2 -fit to various projections

Modern technique: multi-dimensional Maximum Likelihood
fit to $W_{exp}(k_1, \dots, k_n)$

CERNLIB programs : FUMILI, MINUIT

3) Construction of PWA covariant tensor amplitudes

R.E.Behrends ,C.Fronsdal, Phys. Rev. 106, 345 (1957)

A.V. Anisovich et al. J. Phys. G 28 15 (2002)

Refs. specific for baryon-antibaryon final states :

W.H.Liang,P.N.Shen,J.X.Wang, B.S.Zou, J.Phys.G28,333 (2002)

B.S.Zou, F.Hussian, PRC67, 015204 (2003)

J.J.Wu, Z.Ouyang, B.S.Zou, PRC80, 045211 (2009)

Basic ingredients :

(1) spin wave-function for single particle

(2) orbital wave-function for two-particle system

(3) effective couplings

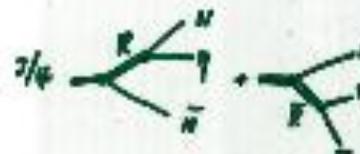
(4) Breit-Wigner propagator, form factor

For ψ decay to baryons:

- Construction of PWA amplitudes.

Rarita-Schwinger covariant tensor formalism.

$$a) \frac{1}{2}^{\pm} N^*$$



$$L_{\text{YME}} = -i g_{\mu\nu} \bar{N} \Gamma R q + \text{h.c.}$$

$$L_{\text{FME}} = -g_{\mu\nu} \bar{R} \Gamma_\mu N \psi^\mu + \frac{i g_{\mu\nu}}{M_N + M_R} \bar{R} \Gamma_\mu \bar{k}_\nu^\mu N \psi^\nu + \text{h.c.}$$

where

$$\Gamma = \Gamma_1, \quad \Gamma_\mu = Y_1 Y_\mu, \quad \Gamma_{\mu\nu} = Y_1 G_{\mu\nu} \quad \text{for } \frac{1}{2}^- N^*,$$

$$\Gamma = \Gamma_2, \quad \Gamma_\mu = Y_2, \quad \Gamma_{\mu\nu} = G_{\mu\nu} \quad \text{for } \frac{1}{2}^+ N^*.$$

$$A_{\lambda^-} = G \bar{q} \left[\frac{K_1 + K_2 + M_{\psi R}}{M_{\psi R}^2 - S_R - i M_{\psi R} \Gamma_{\psi R}} Y_1 Y_\mu E^\mu + Y_2 \not{q} \frac{-K_1 - K_2 + M_{\psi R}}{M_{\psi R}^2 - S_R - i M_{\psi R} \Gamma_{\psi R}} \right] R \\ + G \bar{u} \left[\frac{K_1 + K_2 + M_{\psi R}}{M_{\psi R}^2 - S_R - i M_{\psi R} \Gamma_{\psi R}} Y_1 G_{\mu\nu} \bar{k}_\mu^\nu E^\mu + Y_2 G_{\mu\nu} \not{k}_\mu^\nu \not{q} \frac{-K_1 - K_2 + M_{\psi R}}{M_{\psi R}^2 - S_R - i M_{\psi R} \Gamma_{\psi R}} \right] R$$

$$A_{\lambda^+} = \dots$$

Three basic elements for constructing amplitudes:

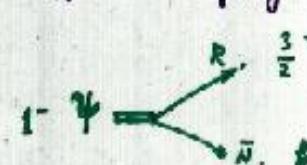
Wave functions, propagators, effective couplings.

$$(1) \quad \frac{3}{2}^- N^*$$

$$\text{Wave function: } U_\lambda(p, s_0) = \sum_{\lambda, s} (i \lambda \frac{1}{2} s / \frac{1}{2} S_0) \partial_\mu(p, \lambda) u(p, s)$$

$$\text{propagators: } P_{\mu\nu} = \frac{p^2 + M_0^2}{p^2 - M_0^2 + i M_0 \Gamma_0} [g_{\mu\nu} - \frac{1}{3} k_\mu k_\nu - \frac{2 p_\mu k_\nu}{3 M_0} + \frac{p_\mu k_\nu - p_\nu k_\mu}{3 M_0}]$$

effective couplings:



$$(1) \bar{R}^\mu \psi^\nu g_{\mu\nu} N$$

$$\bar{R}^\mu \psi^\nu Y_1 k_2 N$$

$$(2) \bar{R}^\mu \psi^\nu Y_2 k_1 N$$

$$\bar{R}^\mu \psi^\nu Y_2 k_2 Y_1 N$$

$$(3) \bar{R}^\mu \psi^\nu k_1 k_2 N$$

$$\bar{R}^\mu \psi^\nu k_1 k_2 Y_1 N$$



$$i \bar{n} \phi Y_1 k_2 R^\mu$$



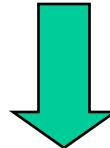
$$i \bar{n} \phi k_2 R^\mu$$

N* with spin J* = n + 1/2

wave function : $u_{\mu_1 \mu_2 \dots \mu_n}(p, n + \frac{1}{2}, s) = \sum_{s_n, s_{n+1}} (n, s_n; \frac{1}{2}, s_{n+1} | n + \frac{1}{2}, s) \epsilon_{\mu_1 \mu_2 \dots \mu_n}(p, n, s_n) u(p, s_{n+1})$

propagator : $G_{R(q)}^{n+\frac{1}{2}(\pm)} = \frac{P_{\mu_1 \mu_2 \dots \mu_n v_1 v_2 \dots v_n}^{n+\frac{1}{2}(\pm)}}{q^2 - m_R^2 + i m_R \Gamma_R}, \quad P_{\mu_1 \mu_2 \dots \mu_n v_1 v_2 \dots v_n}^{n+\frac{1}{2}(\pm)} = \frac{n+1}{2n+3} (\not{p} \pm m) \gamma^\alpha \gamma^\beta P_{\alpha \mu_1 \mu_2 \dots \mu_n \beta v_1 v_2 \dots v_n}^{n+1}$

effective couplings :



FDC (Automatic Feynman Diagram Calculation) – J. X. Wang

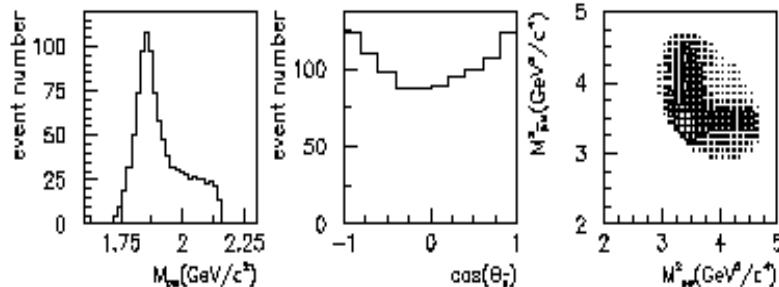


Fortran Programs for amplitudes

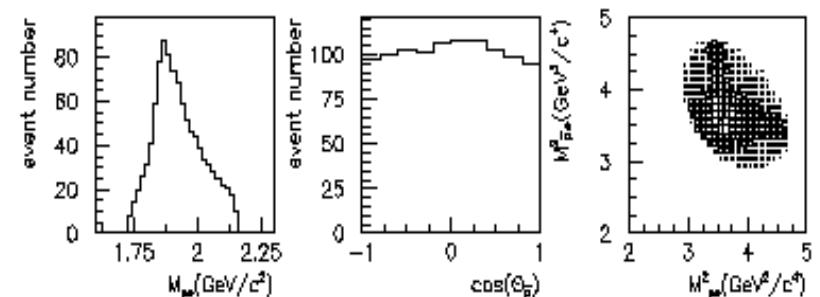


Fit to the data

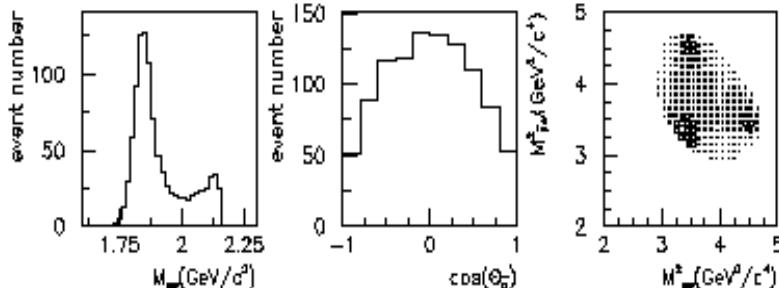
Monte Carlo Simulation for $J/\Psi \rightarrow p + \bar{p} + \omega$



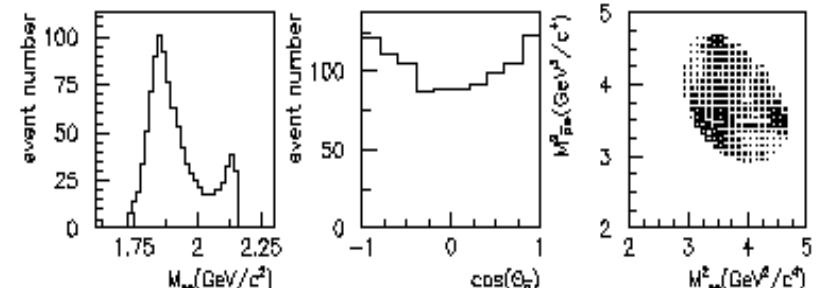
Plots for N^* resonance with $J^P = \frac{1}{2}^-$ and for decay mode 1.



Plots for N^* resonance with $J^P = \frac{1}{2}^+$ and for decay mode 2.

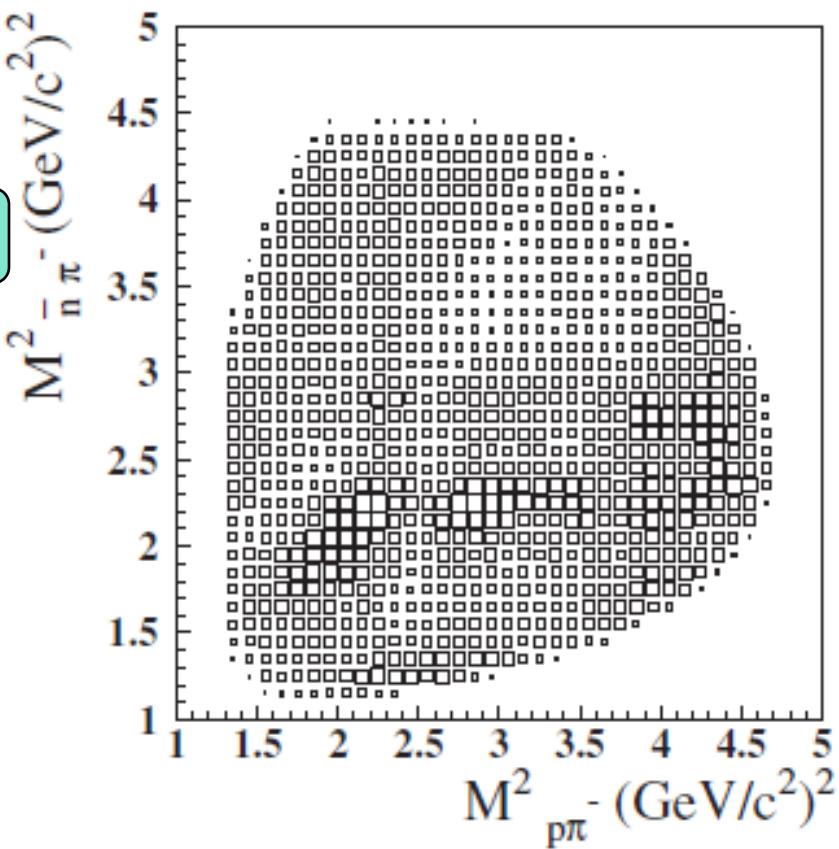
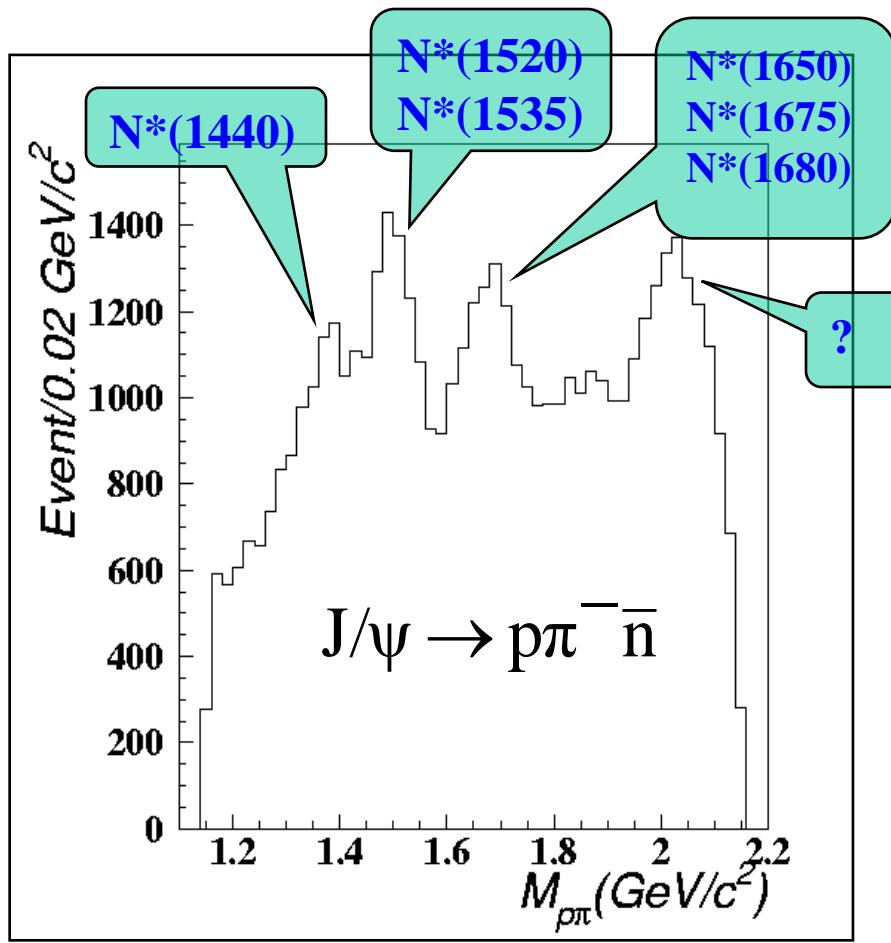


Plots for N^* resonance with $J^P = \frac{1}{2}^+$ and for decay mode 2.



Plots for N^* resonance with $J^P = \frac{5}{2}^-$ and for decay mode 4.

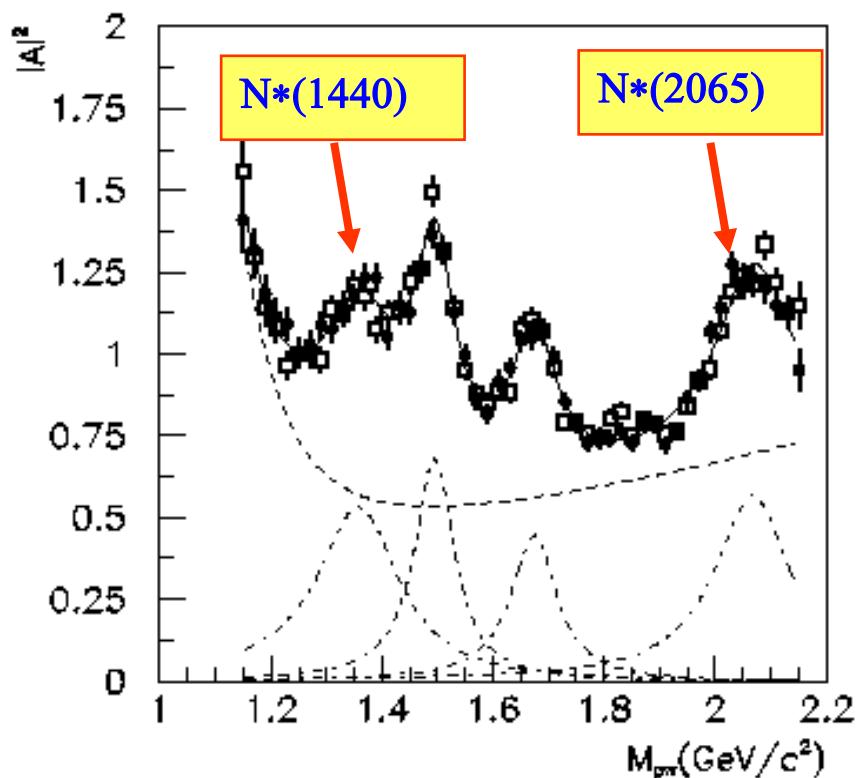
3. N^* resonances observed in $\psi \rightarrow \bar{N} \pi N$



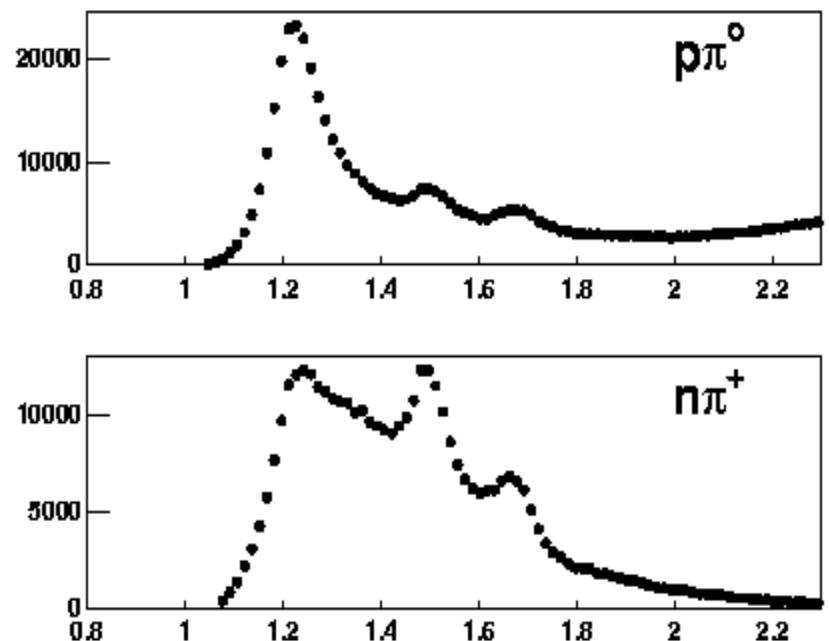
N^* in $\psi \rightarrow \bar{p}n\pi^+$ & $\bar{n}\pi^-p$

$J/\psi \rightarrow \bar{p}n\pi^+$ & $\bar{n}\pi^-p$

V. Burkert



CLAS E=4 GeV ep → eX



Observation of two new N^* peaks in πN mass spectrum

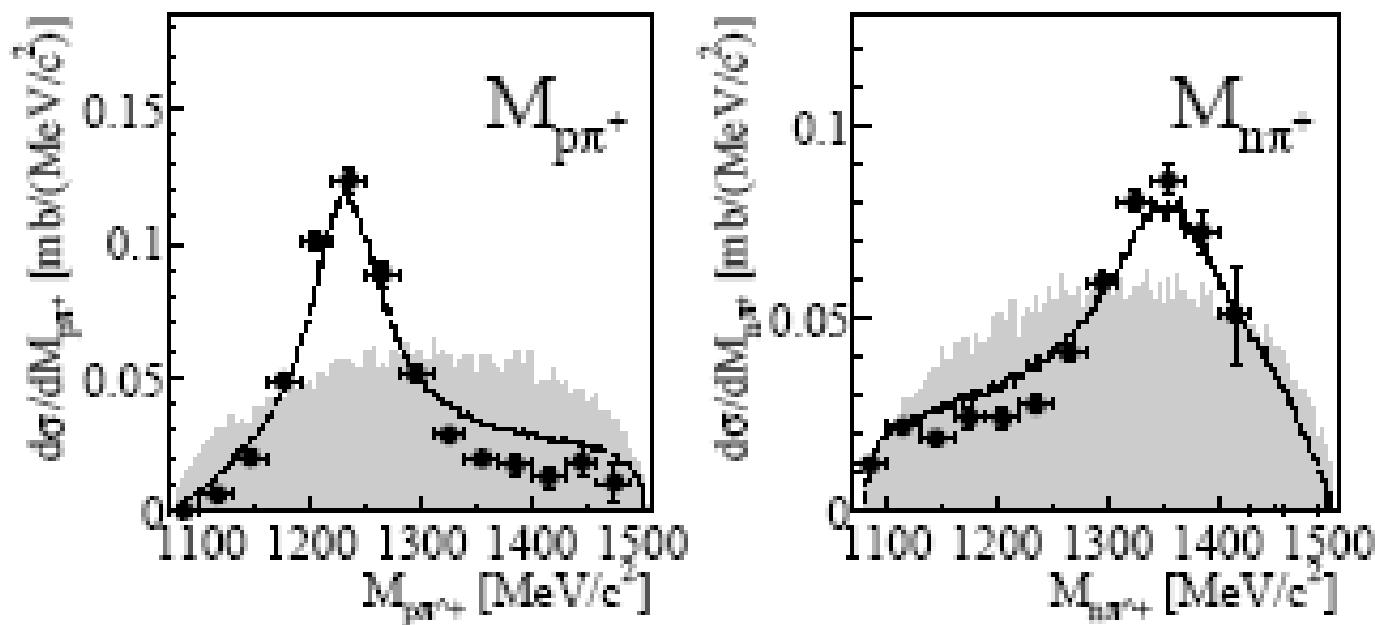
BES, Phys.Rev.Lett.97(2006)062001

The first experiment “seeing” $N^*(1440)$ in πN mass spectrum

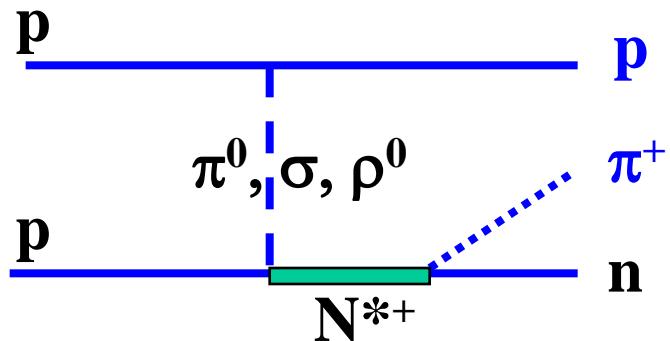
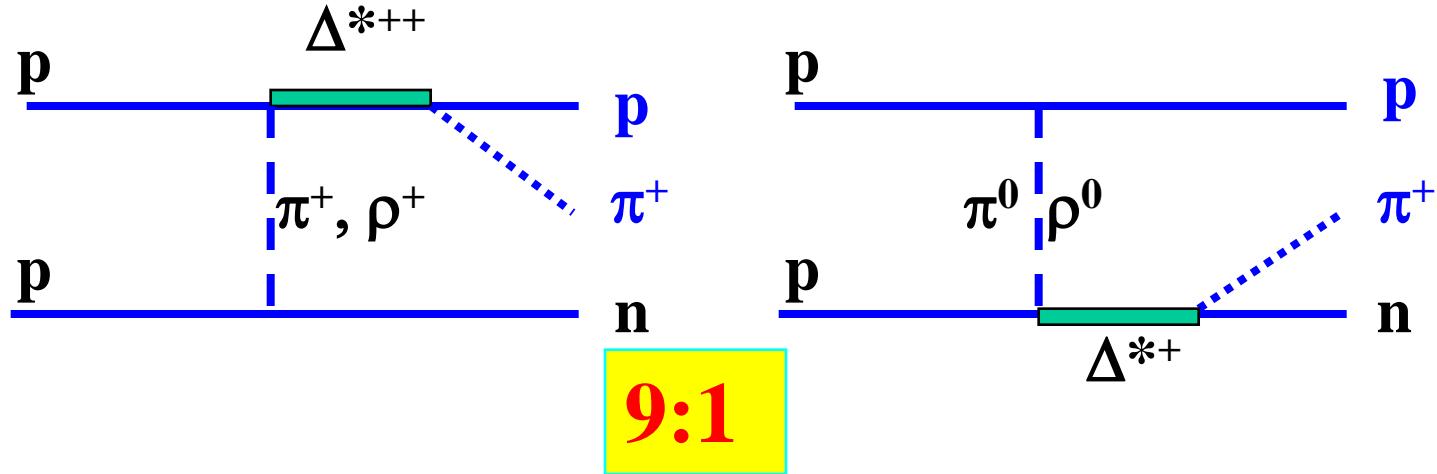
BESII	$M = 1358 \pm 17$,	$\Gamma = 179 \pm 56$	MeV
PDG08	$M = 1365 \pm 15$,	$\Gamma = 190 \pm 30$	MeV

Another experiment “seeing” $N^*(1440)$ in πN mass spectrum

CELSIUS-WASA Collaboration, nucl-ex/0612015



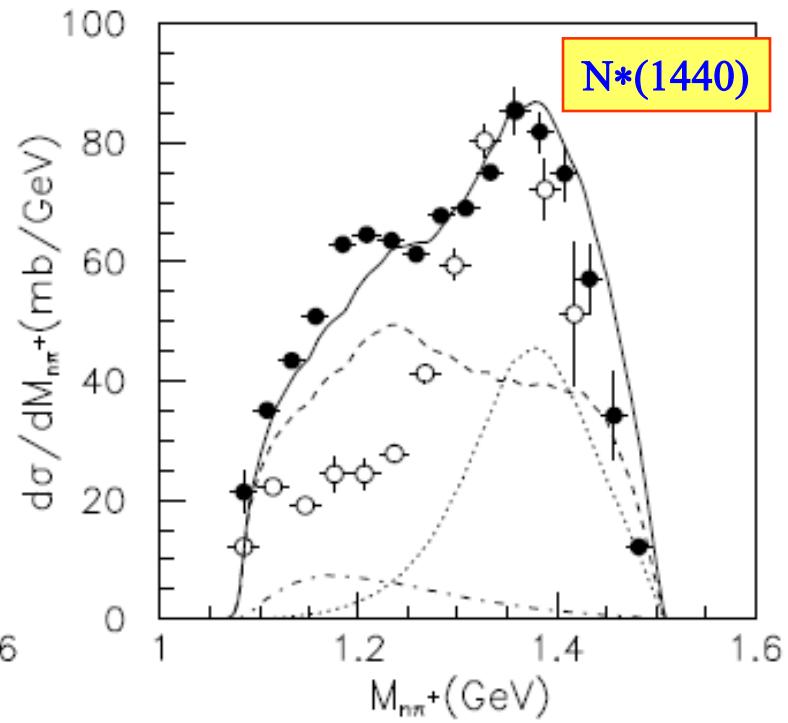
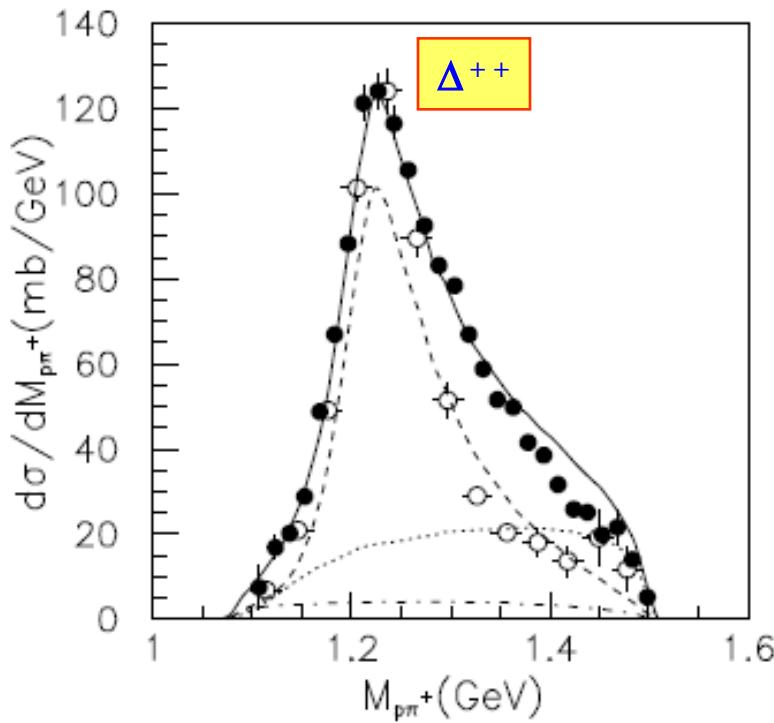
Higher energies and Dalitz plots are needed at COSY&CSR !



A quite good isospin filter !
Look for “missing” N^* by
 σp production !

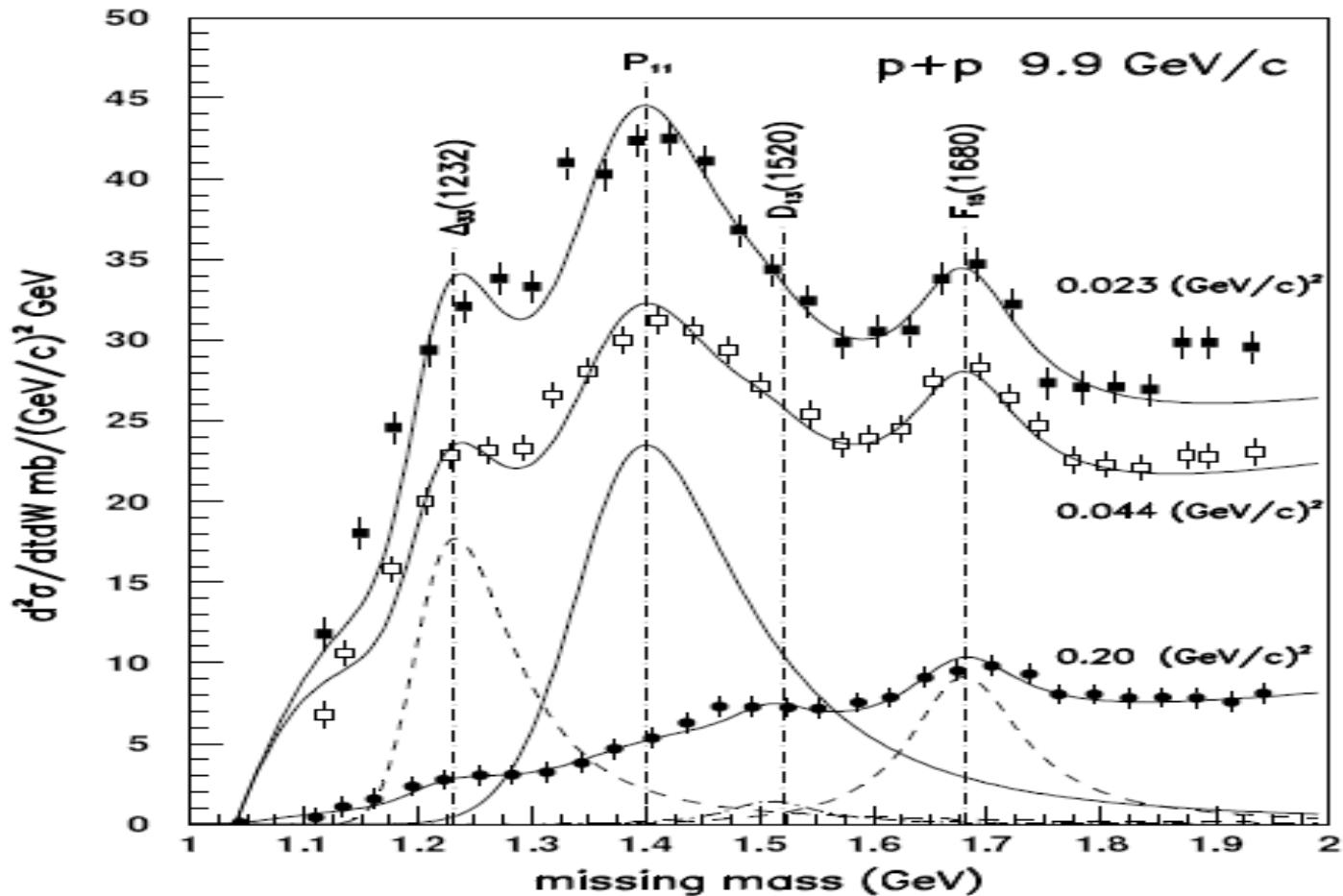
Theoretical study on $pp \rightarrow pn\pi^+$

Z.Ouyang,J.J.Xie,B.S.Zou &H.S.Xu, Nucl.Phys.A 821(2009)220; IJMPE(2009)



Including $N^*(1440)$ by σ exchange is crucial to reproduce data

N*(1440) in missing mass of $\alpha + p \rightarrow \alpha + X$ & $p + p \rightarrow p + X$



$p + p \rightarrow p + X$



A.V. Sarantsev et al., Phys. Lett. B659, 94 (2008).



X.Cao, B.S.Zou, H.S.Xu, Phys. Rev. C 81(2010) 06***

Both analyses demand larger $\sigma N/\pi\Delta$ ratio for $N^*(1440)$

The new $N^*(2065)$ peak in πN mass spectrum

$$M = 2068 \pm 3 {}^{+15}_{-40} \text{ MeV}/c^2, \Gamma = 165 \pm 14 \pm 40 \text{ MeV}/c^2$$

$J/\psi \rightarrow \bar{n} N^*(2065)$ with $L=0$ (small excess energy)

limits its J^P to be $3/2^+$ or $1/2^+$ (spin filter)

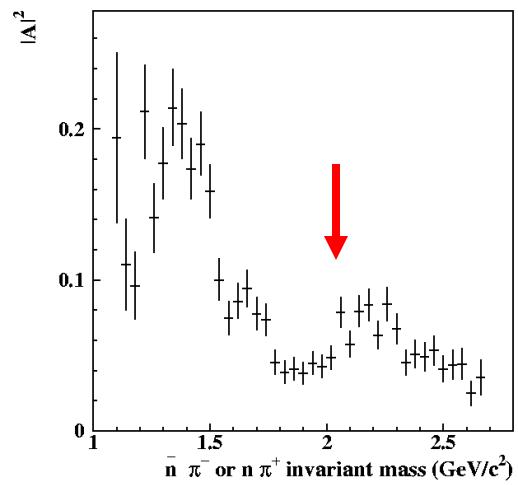
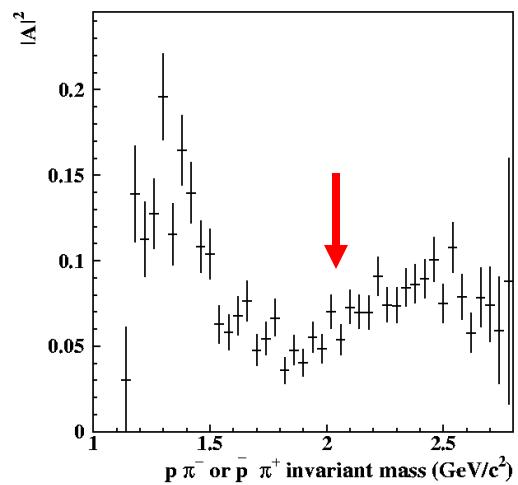
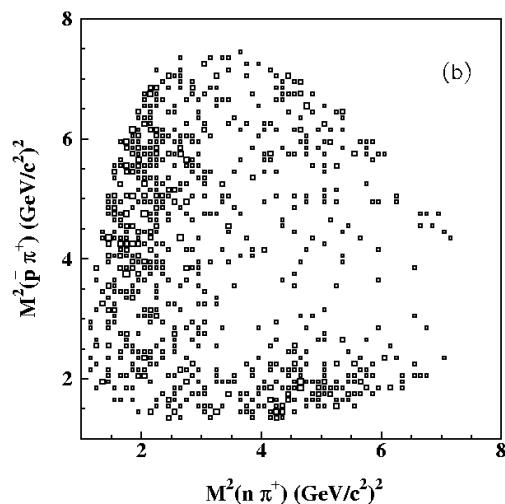
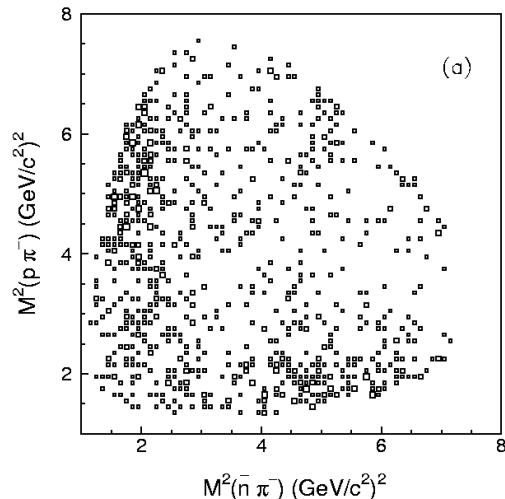
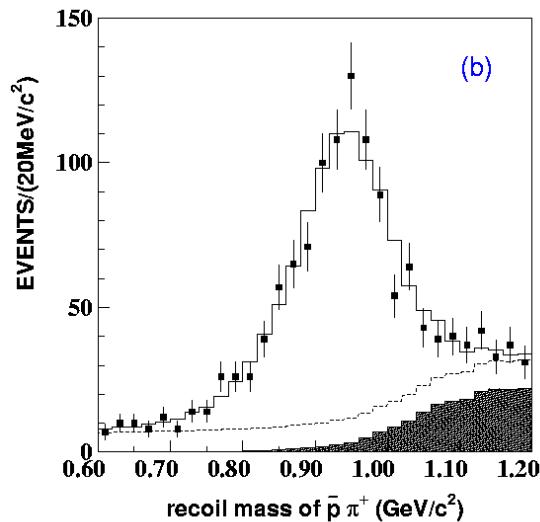
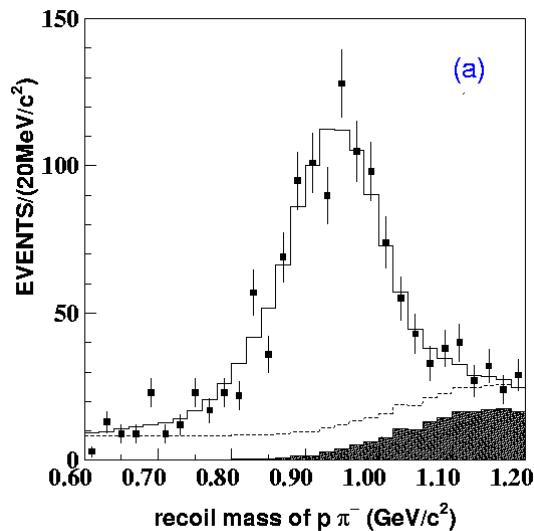
PWA with event-based Maximum-Likelihood method

→ Both $3/2^+$ and $1/2^+$ are there.

$1/2^+$ or $3/2^+$ (improve log likelihood by 400)

$1/2^+ + 3/2^+$ (improve log likelihood further by 60)

$$\psi(2S) \rightarrow p\pi^-\bar{n} + c.c.$$



$$\gamma p \rightarrow \pi^+ n \text{ & } \gamma n \rightarrow \pi^- p$$

Jlab Hall A Collaboration,
L.Y.Zhu et al., PRL91 (2003) 022003

**Clear broad resonance
cluster around 2.2 GeV**

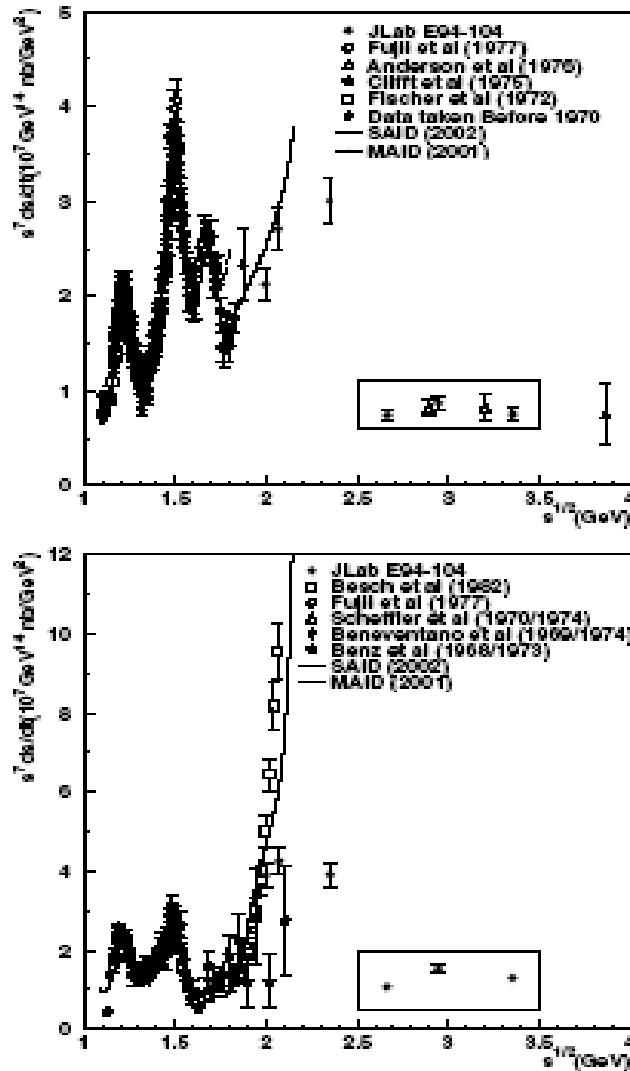
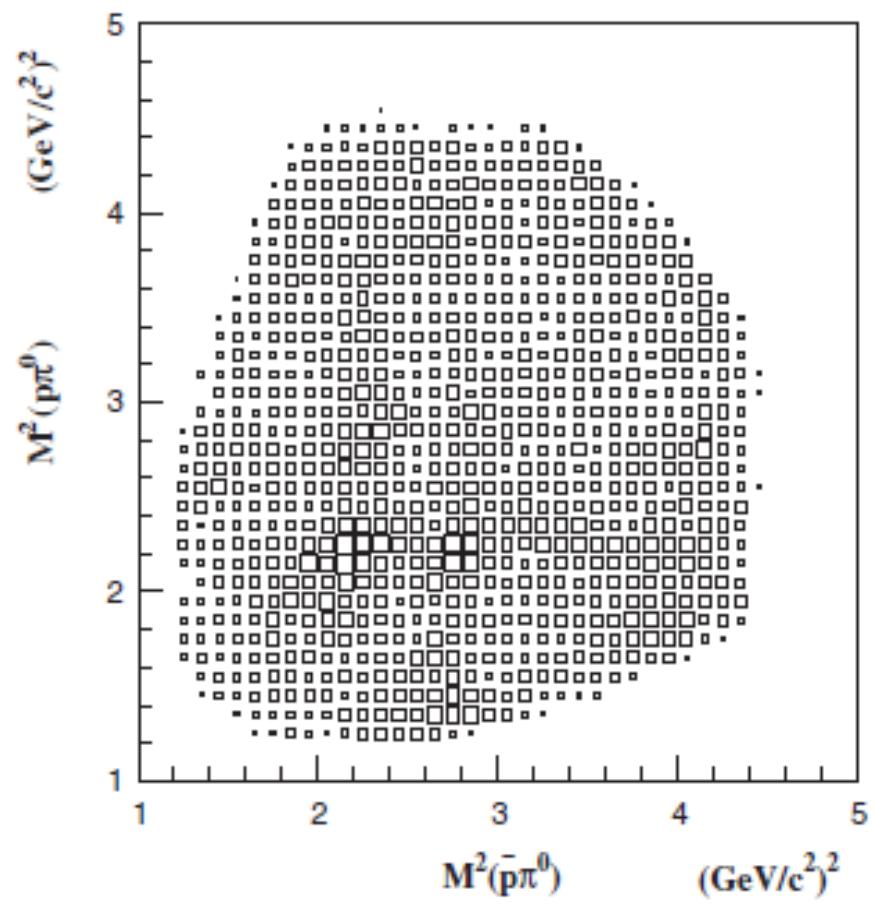
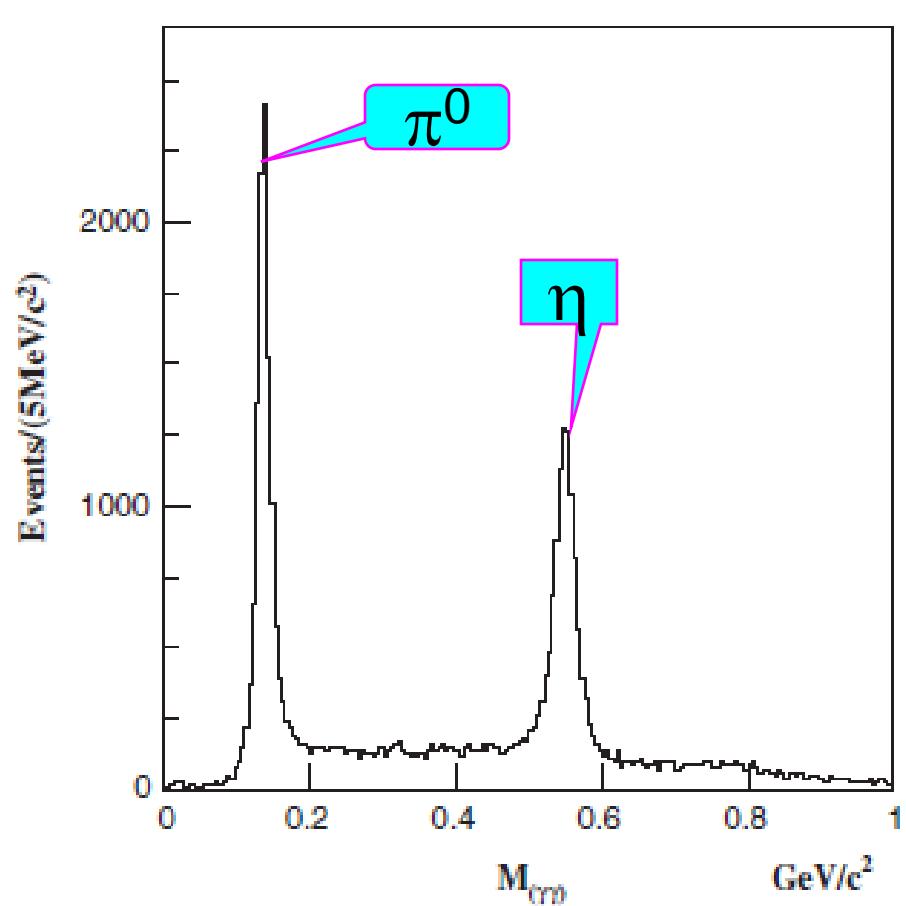
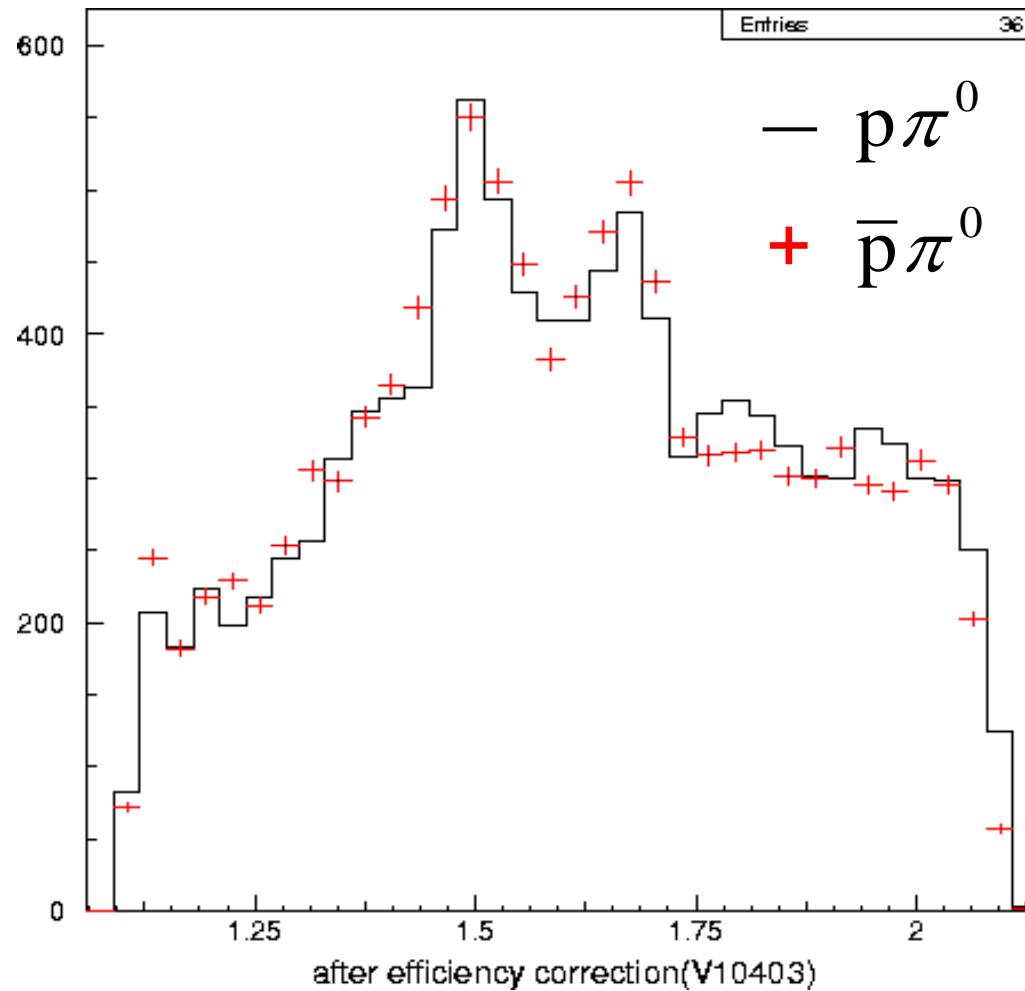


FIG. 1. The scaled differential cross section $s^7 \frac{d\sigma}{dt}$ versus center-of-mass energy for the $\gamma p \rightarrow \pi^+ n$ (upper plot) and $\gamma n \rightarrow \pi^- p$ (lower plot) at $\theta_{c.m.} = 90^\circ$. The data from JLab E94-104 are shown as solid circles. The error bars for the new

N^* in $J/\psi \rightarrow p\bar{p}\pi^0$



$N^* \text{ in } J/\psi \rightarrow p\bar{p}\pi^0$



N^* in $J/\psi \rightarrow p\bar{p}\pi^0$

- More than 10 K events were selected, and detailed PWA was performed
- $N(1440)$, $N(1520)$, $N(1535)$, $N(1650)$, $N(1675)$, $N(1680)$, $N(1710)$ are needed.
- $N(2065)$ is needed in the fit, and J^{PC} favor $3/2+$

$$M = 2040 {}^{+3}_{-4} \pm 25 \text{ MeV}/c^2, \Gamma = 230 \pm 8 \pm 56 \text{ MeV}/c^2$$

Resonance	Mass (MeV/c ²)	(Width MeV/c ²)	J ^P	Fraction (%)
N(1440)	1455 ⁺² ₋₇	316 ⁺⁵ ₋₆	1 ⁺ 2 ⁻	16.37
N(1520)	1513 ⁺³ ₋₄	127 ⁺⁷ ₋₈	3 ⁻ 2 ⁻	7.96
N(1535)	1537 ⁺² ₋₆	135 ⁺⁸ ₋₈	1 ⁻ 2 ⁻	7.58
N(1650)	1650 ⁺³ ₋₆	145 ⁺⁵ ₋₁₀	1 ⁻ 2 ⁻	9.06
N(1710)	1715 ⁺² ₋₂	95 ⁺² ₋₁	1 ⁺ 2 ⁻	25.33
N _x (2065)	2040 ⁺³ ₋₄	230 ⁺⁸ ₋₈	3 ⁺ 2 ⁻	23.39

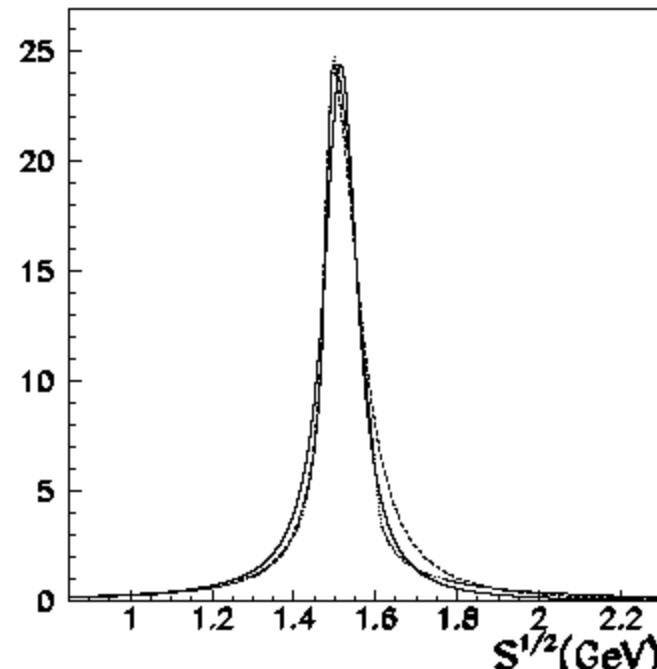
$$\Gamma_{N(1440)(s)} = \Gamma_{N(1440)} (0.7 \frac{B_1(q_{\pi N}) \rho_{\pi N}(s)}{B_1(q_{\pi N}^{N^*}) \rho_{\pi N}(M_{N^*}^2)} + 0.3 \frac{B_1(q_{\pi \Delta}) \rho_{\pi \Delta}(s)}{B_1(q_{\pi \Delta}^{N^*}) \rho_{\pi \Delta}(M_{N^*}^2)})$$

Mass of $N^*(1535)$

$$BW(p_{N^*}) = \frac{1}{M_{N^*}^2 - s - iM_{N^*}\Gamma_{N^*}(s)}$$

(1) $\Gamma_{N^*}(s) = 98\text{MeV}$

$$M_{N^*} = 1515\text{MeV}$$

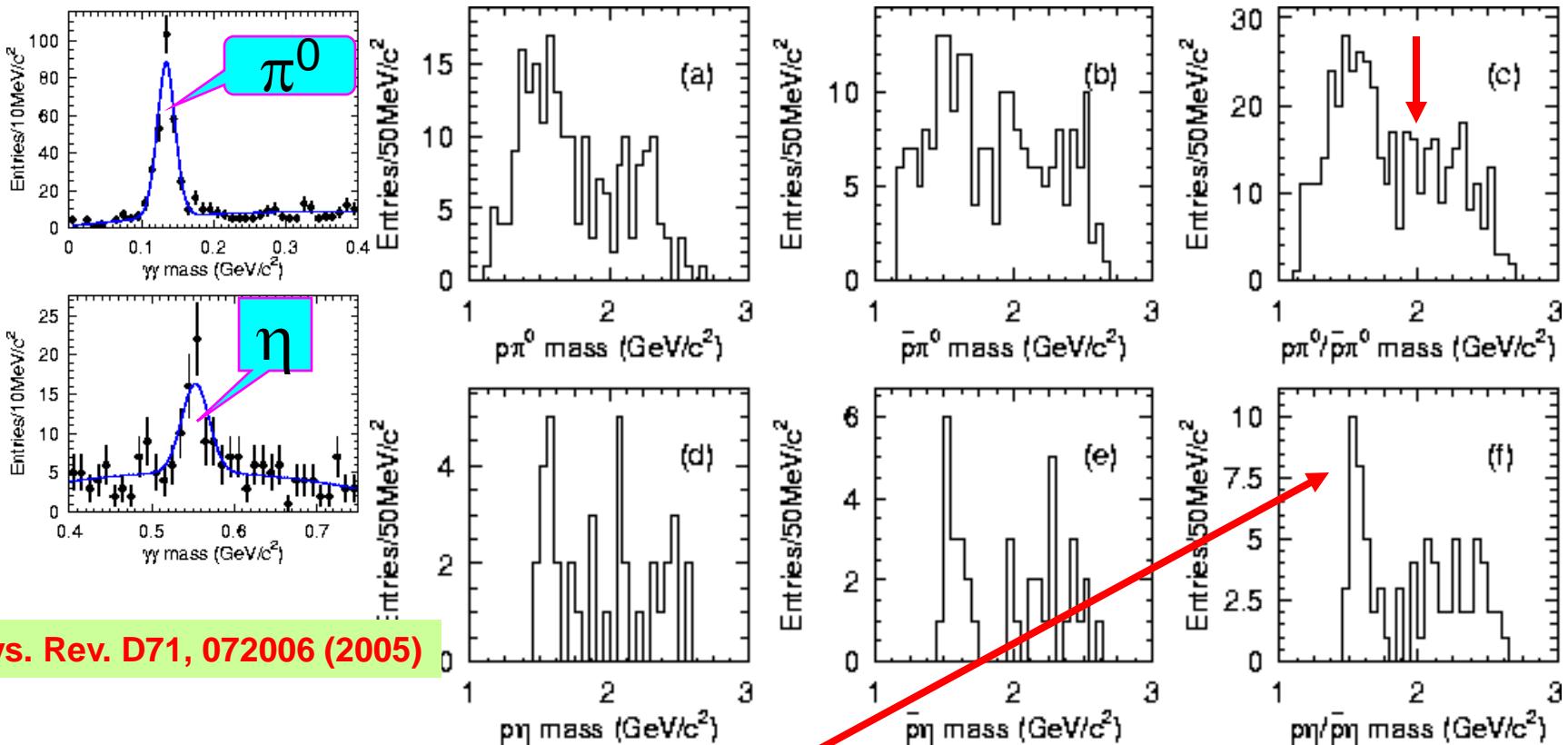


(2) $\Gamma_{N^*}(s) = \Gamma_{N^*}^0 \left(0.5 \frac{\rho_{\pi N}(s)}{\rho_{\pi N}(M_{N^*}^2)} + 0.5 \frac{\rho_{\eta N}(s)}{\rho_{\eta N}(M_{N^*}^2)} \right) = \Gamma_{N^*}^0 [0.8\rho_{\pi N}(s) + 2.1\rho_{\eta N}(s)]$

$$M_{N^*} = 1535 \text{ MeV} \text{ and } \Gamma_{N^*}^0 = 150 \text{ MeV}$$

(3) $\Gamma_{N^*}(s) = \Gamma_{N^*}^0 [0.8\rho_{\pi N}(s) + 2.1\rho_{\eta N}(s) + 3.5\rho_{\Lambda K}(s)]$ $M_{N^*} \approx 1400\text{MeV}$
 $\Gamma_{N^*}^0 = 270\text{MeV}$

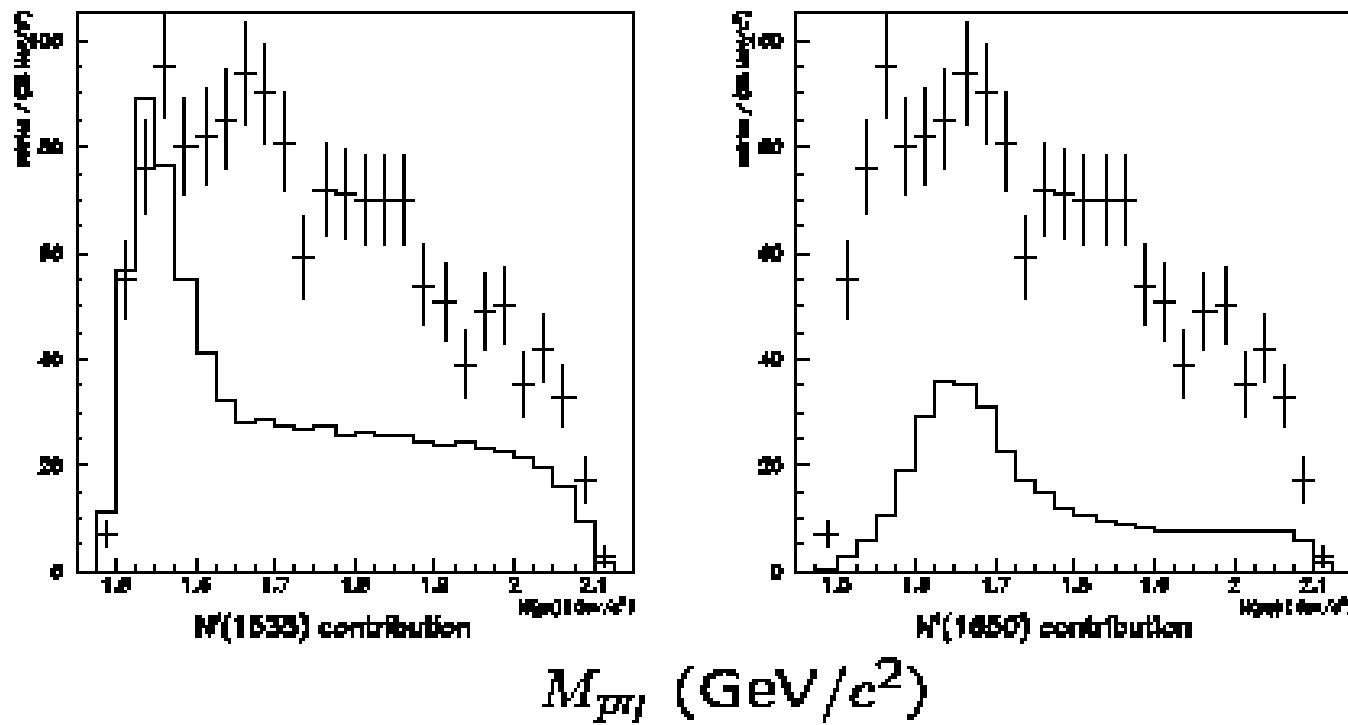
$\psi(2S) \rightarrow p\bar{p}\pi^0(\eta)$



Phys. Rev. D71, 072006 (2005)

- A faint accumulation of events around $2065 \text{ MeV}/c^2$
- The enhancement between 1.4 and $1.7 \text{ GeV}/c^2$
- Possible $N^*(1535)$ in the $p\eta$ invariant mass

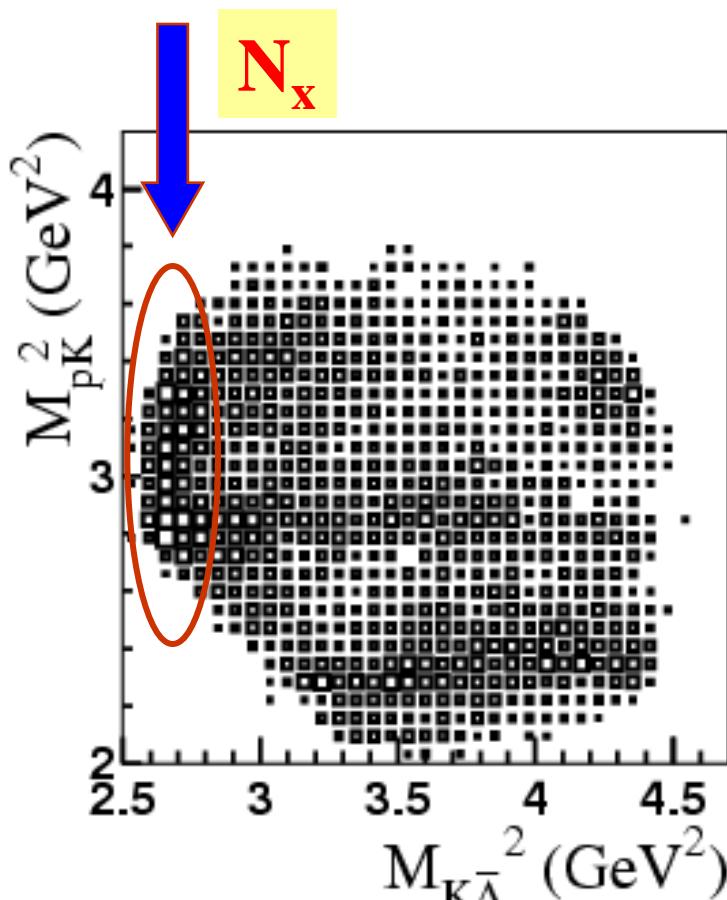
PWA Results from $J/\psi \rightarrow p\bar{p}\eta$ (BES I 7.8M)



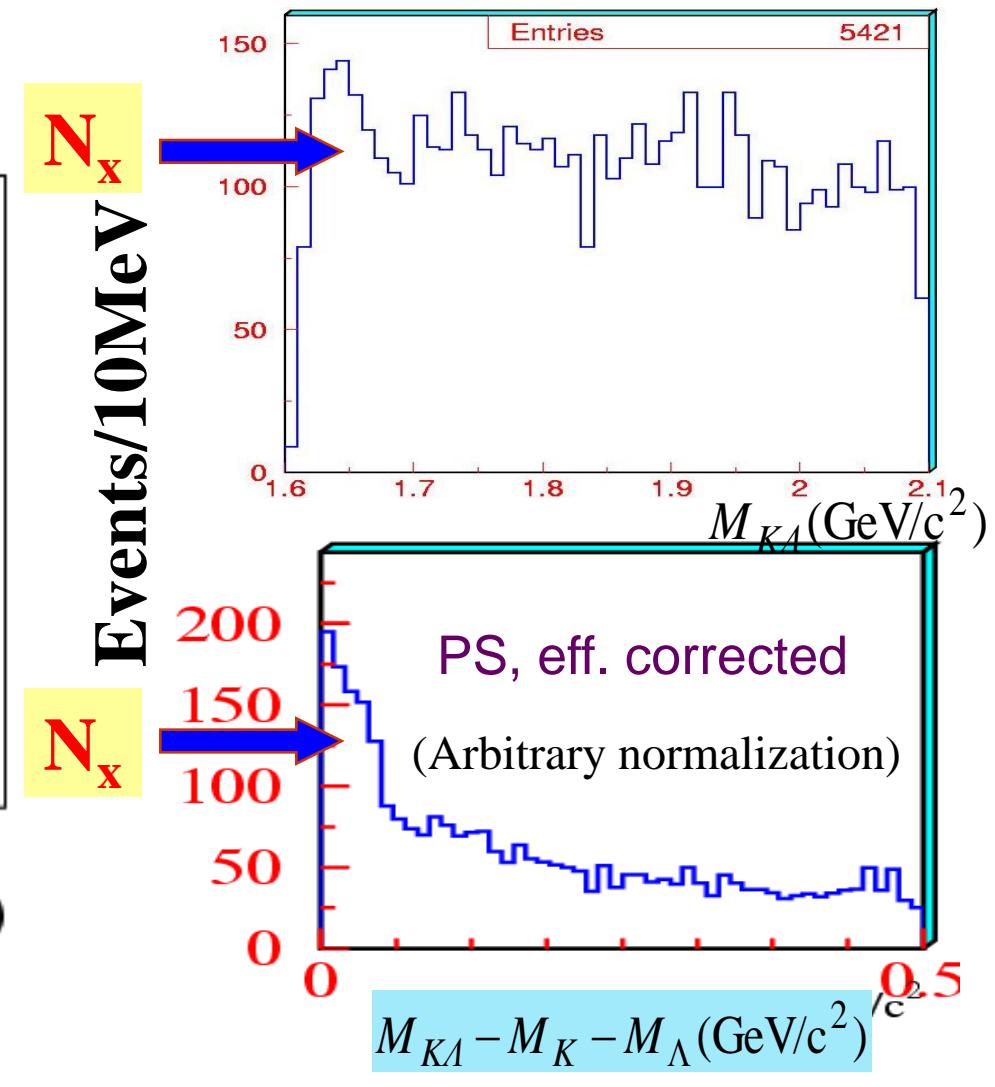
Phys. Lett. B510 (2001) 75

4. N^* and Λ^* observed in $J/\psi \rightarrow \bar{\Lambda} K N$

N^* in $J/\psi \rightarrow p K^- \bar{\Lambda} + c.c.$

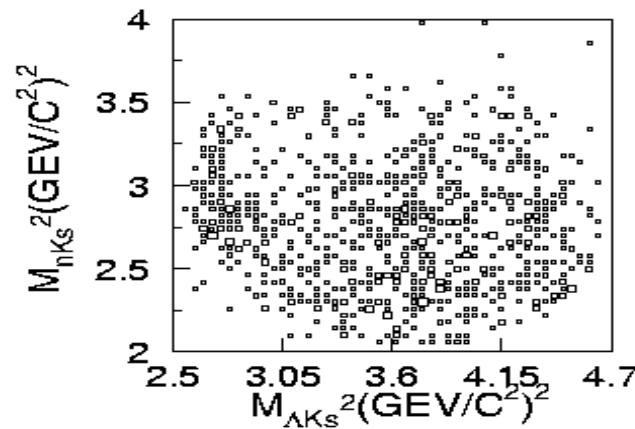
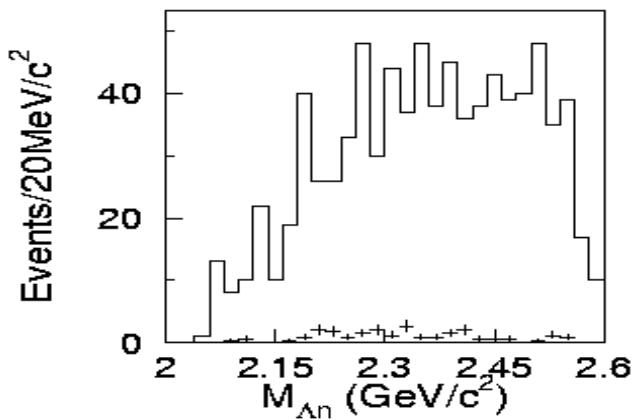
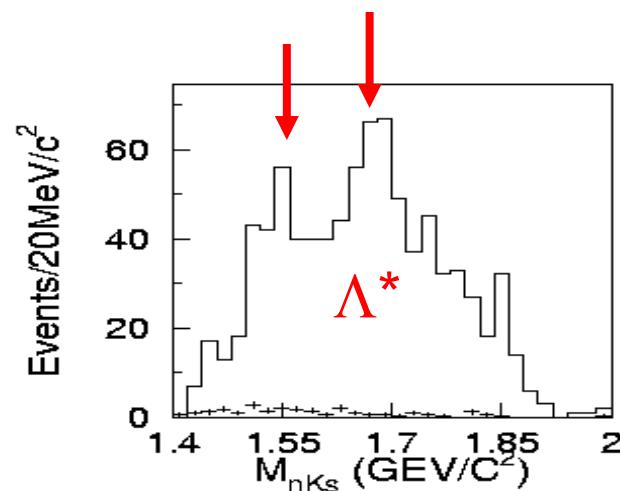
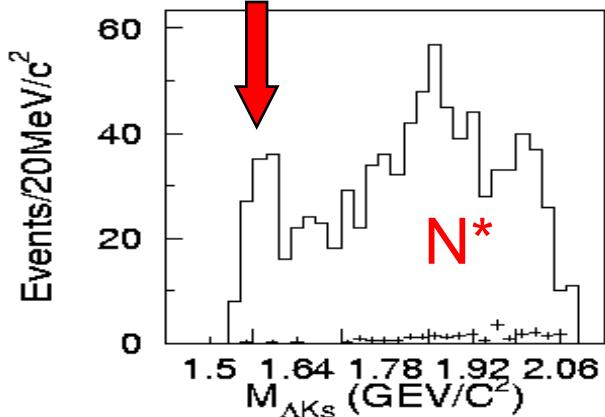


Mass **1500~1650MeV**
Width **70~110MeV**
 J^P **favors 1/2-**



$J/\psi \rightarrow nK_S^0 \bar{\Lambda}$

Phys. Lett. B659 (2008) 789



- An enhancement near ΛK_S threshold is evident
- N^* and Λ^* found in the ΛK_S and $n K_S$ spectrum

a) Assuming N_x to be purely $N^*(1535)$:

B.C. Liu, B.S. Zou, PRL96 (2006) 042002; PRL98 (2007) 039102

From relative branching ratios of
 $J/\psi \rightarrow p \bar{N}^* \rightarrow p (K^- \bar{\Lambda}) / p (\bar{p}\eta)$



$$g_{N^*K\Lambda} / g_{N^*p\eta} / g_{N^*N\pi} \sim 2 : 2 : 1$$

b) N_x as dynamical generated with unitary chiral theory:

$N^*(1535)$ + non-resonant part

L.S.Geng, E.Oset, B.S. Zou, M.Doring, PRC79 (2009) 025203

$$g_{N^*K\Lambda} / g_{N^*p\eta} / g_{N^*N\pi} \sim 1.2 : 2 : 1$$

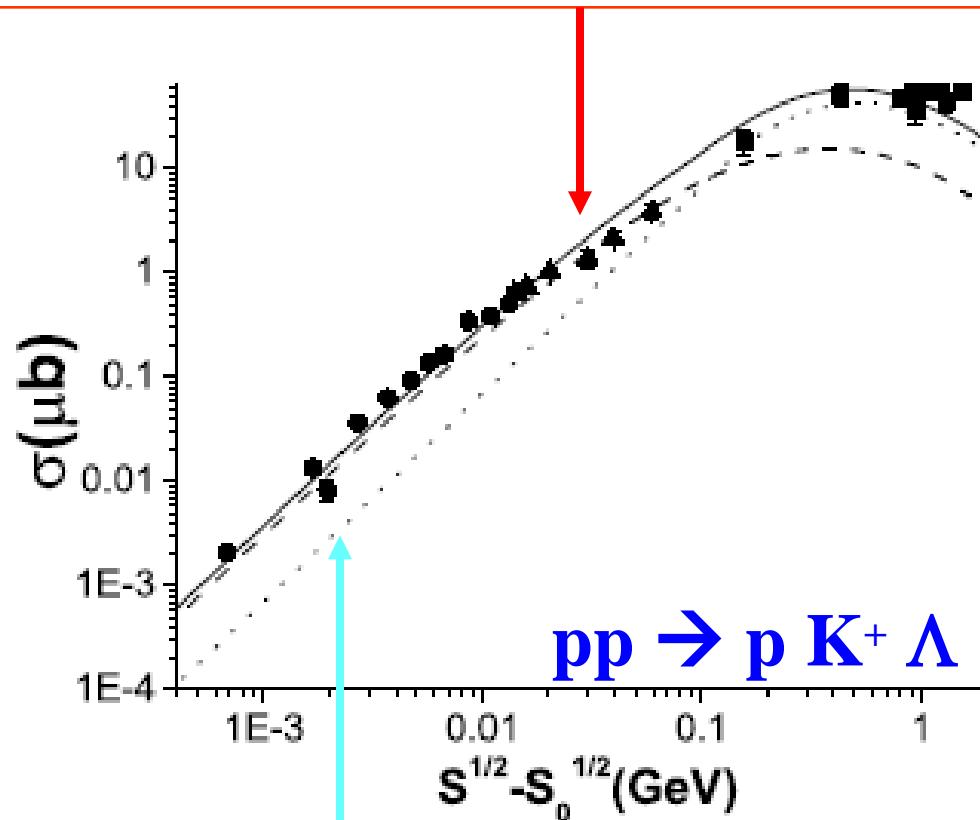
Phenomenology : Large $g_{N^*K\Lambda}$ \rightarrow large $\bar{s}s$ in $N^*(1535)$

$\bar{s}[su][ud]$ or $K\Lambda-K\Sigma$ state

Evidence for large $g_{N^* K\Lambda}$ from $pp \rightarrow p K^+ \Lambda$

Total cross section and theoretical results with
 $N^*(1535)$, $N^*(1650)$, $N^*(1710)$, $N^*(1720)$

B.C.Liu, B.S.Zou, Phys. Rev. Lett. 96 (2006) 042002



Tsushima,Sibirtsev,Thomas, PRC59 (1999) 369, without including $N^*(1535)$

FSI vs N*(1535) contribution in $pp \rightarrow p K^+ \Lambda$

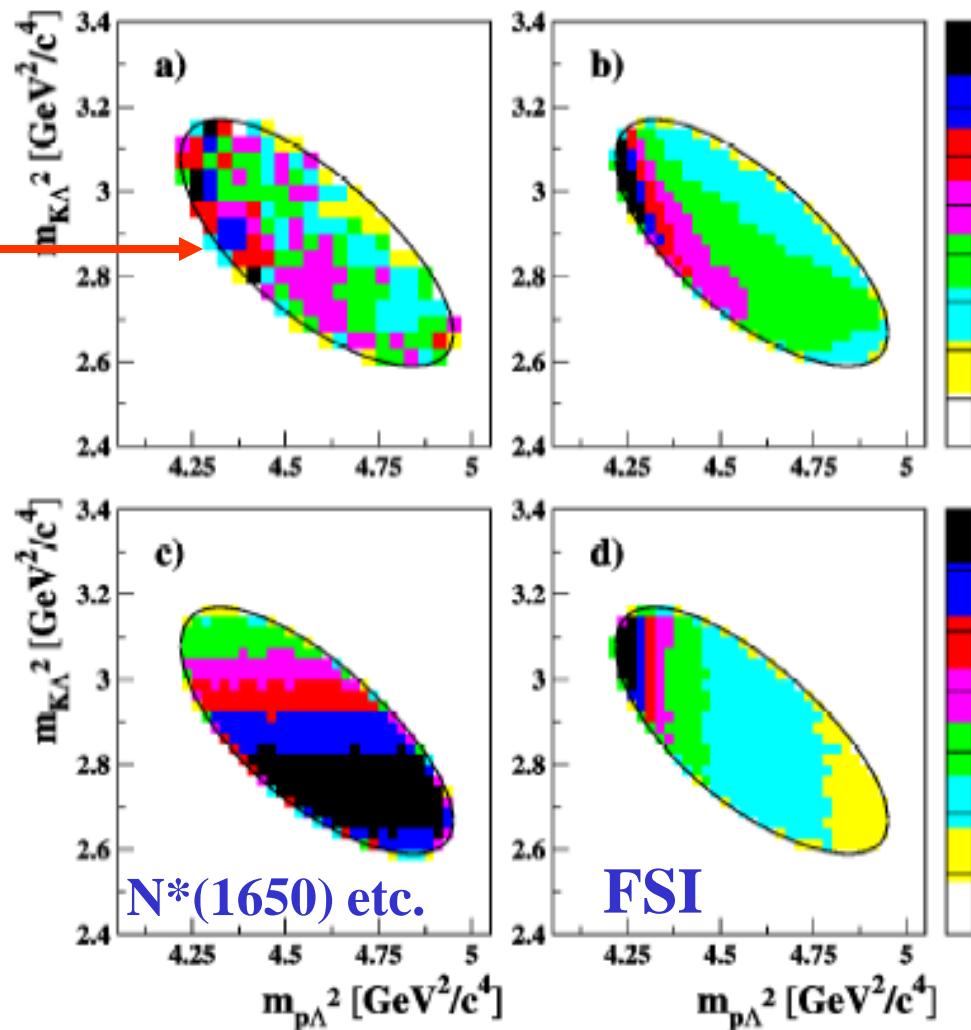
B.C.Liu & B.S.Zou, Phys. Rev. Lett. 98 (2007) 039102 (reply)

A.Sibirtsev et al., Phys. Rev. Lett. 98 (2007) 039101 (comment)

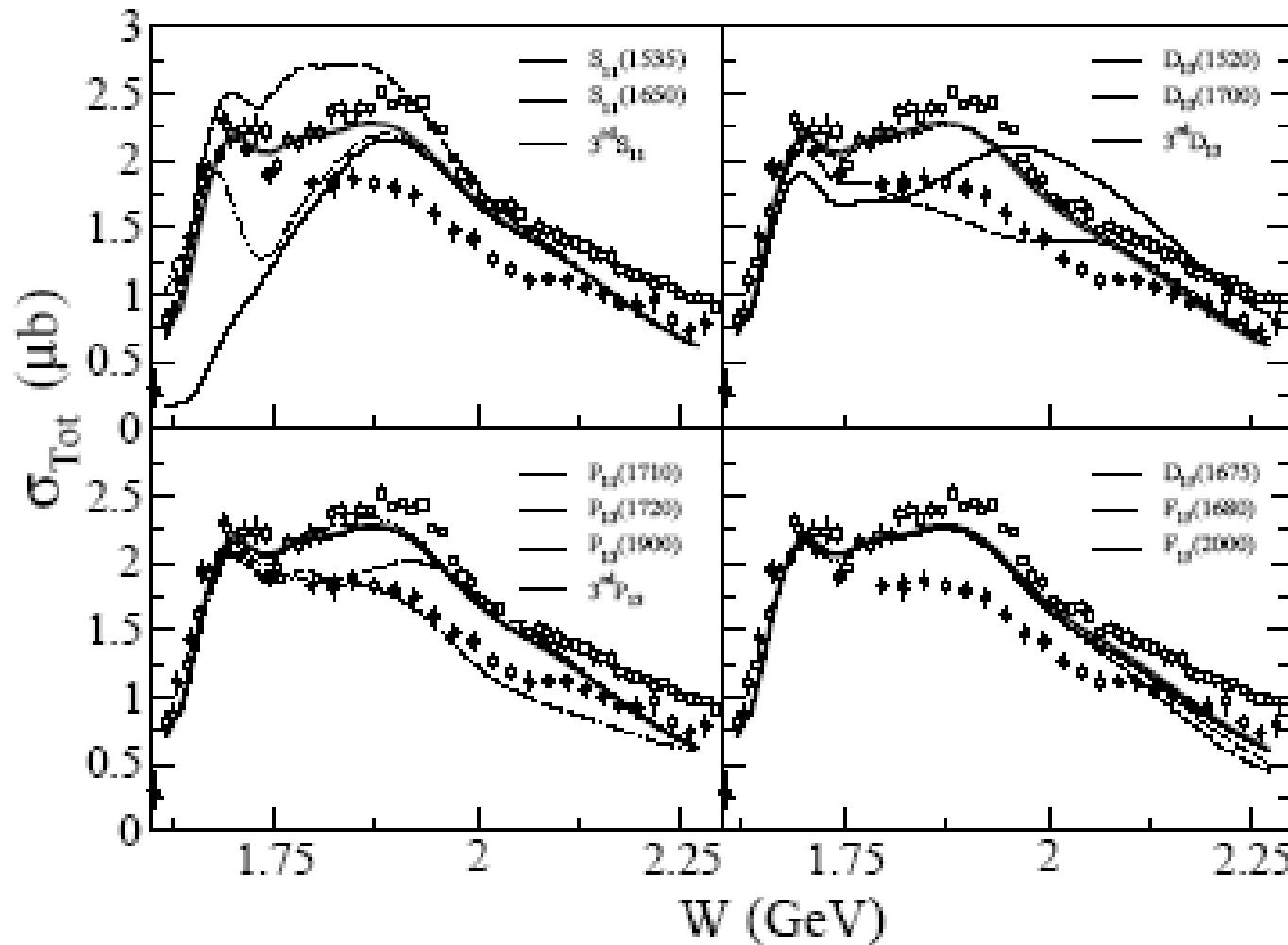
COSY-TOF data
S. Abdel-Samad *et al.*,
Phys.Lett.B632:27(2006)



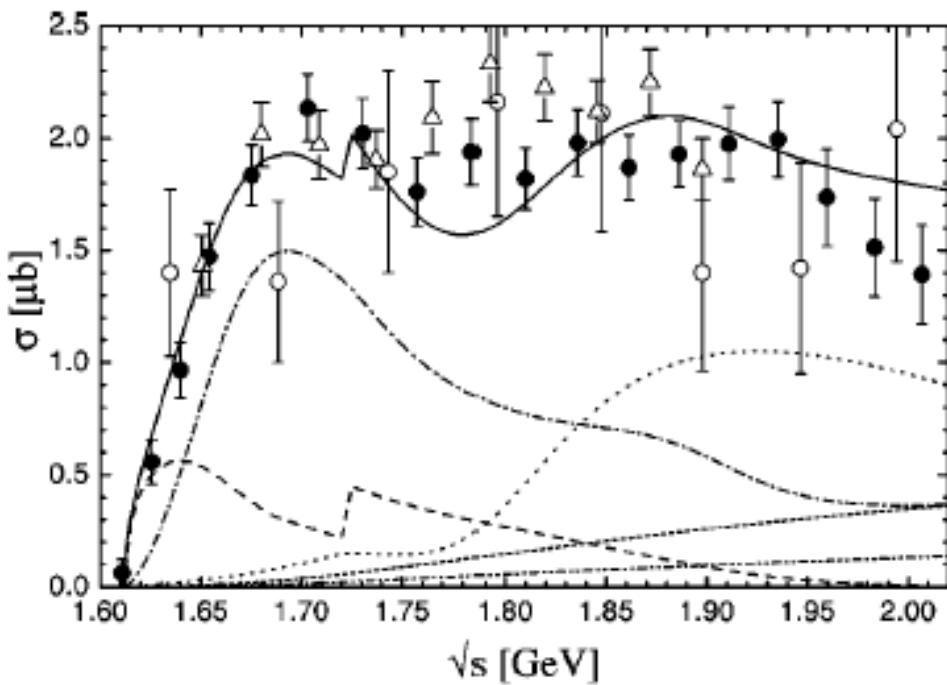
**Both FSI & N*(1535)
are needed !**



Evidence for large $g_{N^* K\Lambda}$ from $\gamma p \rightarrow K^+ \Lambda$



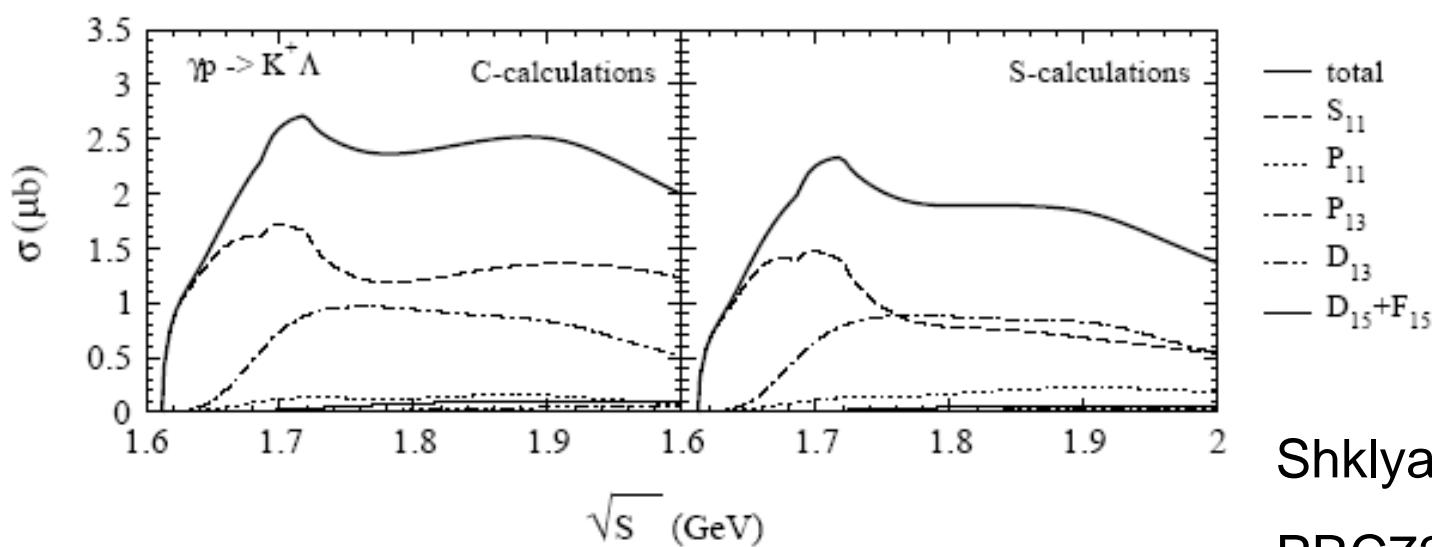
B. Julia-Diaz, B. Saghai, T.-S.H. Lee, F. Tabakin, Phys. Rev. C 73, 055204 (2006)



G.Penner&U.Mosel,
PRC66 (2002) 055212

Partial wave decomposition
For the fit to SAPHIR92-94
Data

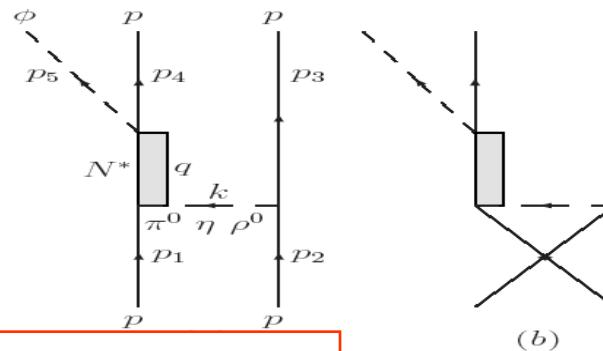
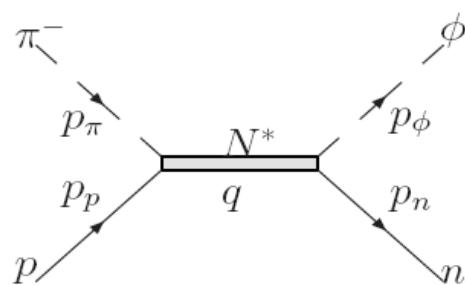
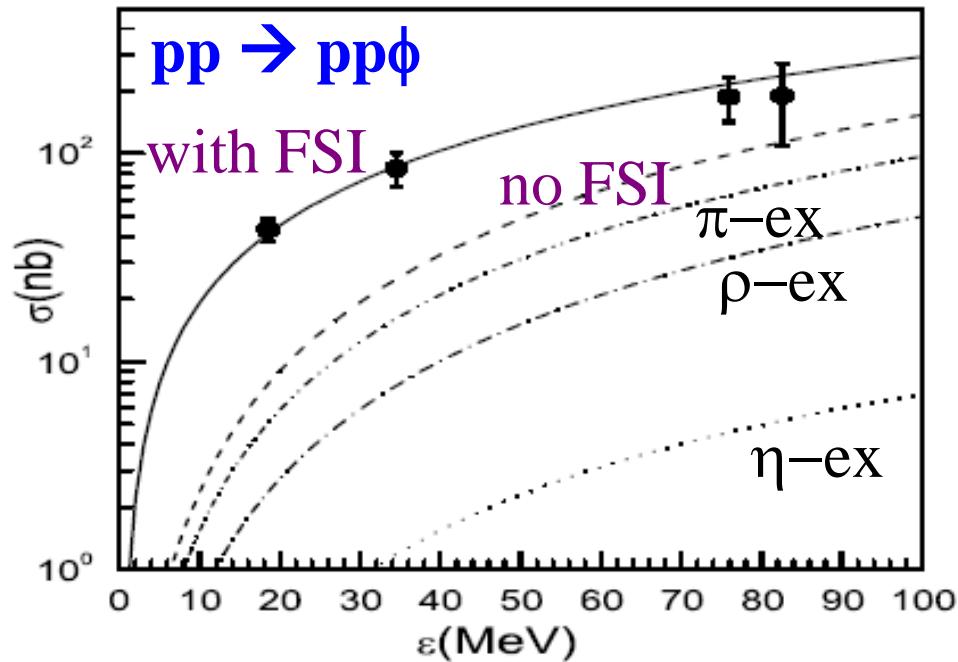
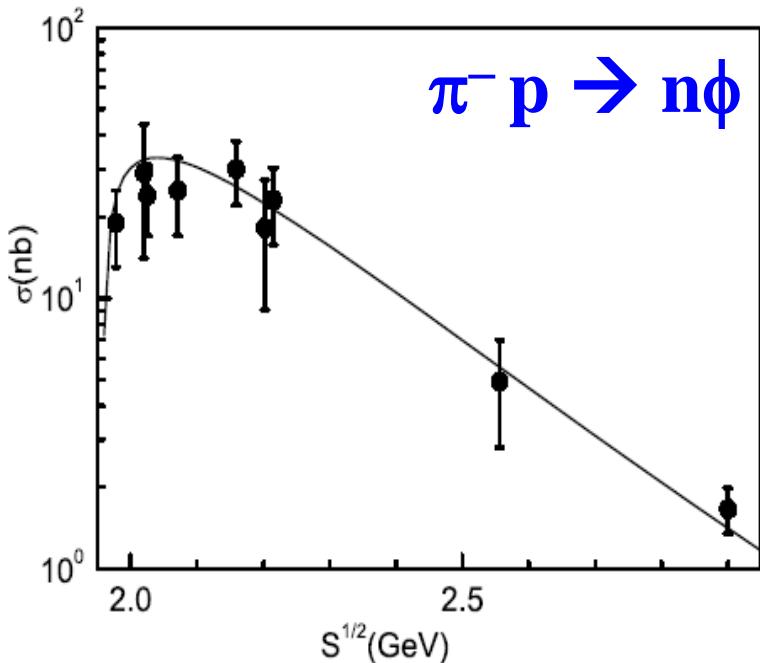
Dashed line : 1/2 -
Dot-dashed line : 3/2+



Shklyar,Lenske&Mosel,
PRC72 (2005) 015210

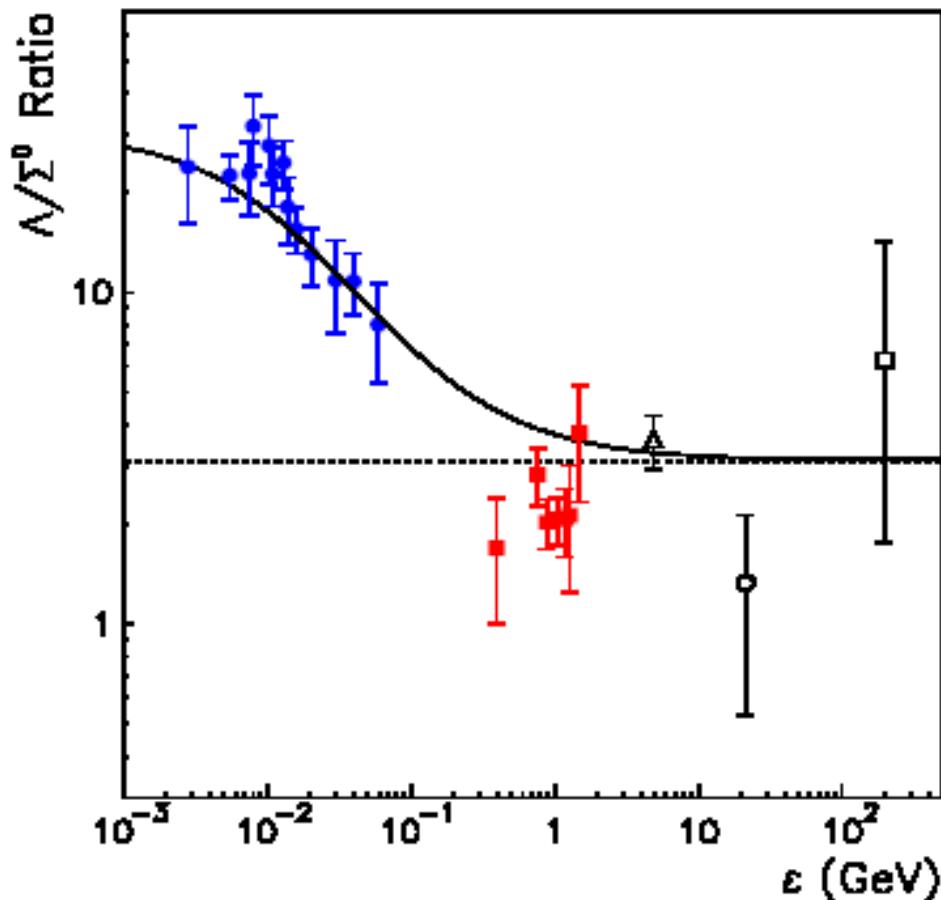
Evidence for large $g_{N^*N\phi}$ from $\pi^- p \rightarrow n\phi$ & $pp \rightarrow pp\phi$

Xie, Zou & Chiang, arXiv:0705.3950



Evasion of OZI rule by $N^*(1535)$!

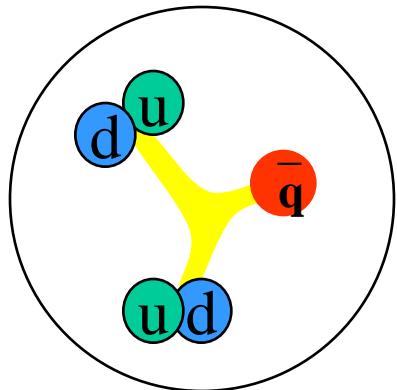
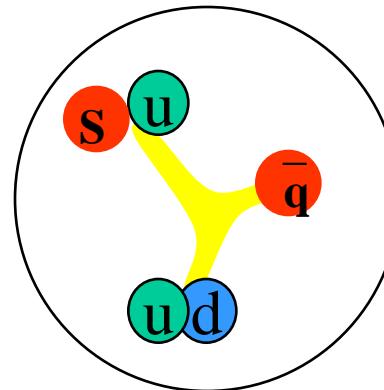
Evidence for weaker $g_{N^*K\Sigma}$ from $pp \rightarrow p K^+ \Lambda$ / $pp \rightarrow p K^+ \Sigma^0$



A.Sibirtsev et al.,
EPJA29 (2006) 363

Fig. 3. The Λ/Σ^0 cross-section ratio as a function of the excess energy ϵ . The solid circles show the ratio obtained for the $pp \rightarrow K^+ \Lambda p$ and $pp \rightarrow K^+ \Sigma^0 p$ reactions at COSY [2]. Solid

Nature of $\mathbf{N}^*(1535)$ and its $1/2^-$ octet partner


$$\begin{array}{c} \bar{q} \\ [ud] \\ [ud] \end{array} \}^{1/2+}_{L=1}$$

$$\begin{array}{c} \bar{q} \\ [ud] \\ [us] \end{array} \}^{1/2-}_{L=0}$$

Zhang et al, hep-ph/0403210

$$N^*(1535) \sim uud \ (L=1) + \varepsilon [ud][us] \bar{s} + \dots$$

$$N^*(1440) \sim uud \ (n=1) + \xi [ud][ud] \bar{d} + \dots$$

$$\Lambda^*(1405) \sim uds \ (L=1) + \varepsilon [ud][su] \bar{u} + \dots$$

$N^*(1535)$: $[ud][us] \bar{s} \rightarrow$ larger coupling to $N\eta, N\eta', N\phi$ & $K\Lambda$, weaker to $N\pi$ & $K\Sigma$, and heavier !

The new picture for the $1/2^-$ octet predicts:

Λ^* [us][ds] \bar{s} ~ 1575 MeV

Σ^* [us][du] \bar{d} ~ 1360 MeV

Ξ^* [us][ds] \bar{u} ~ 1520 MeV

J/ψ decay branching ratio * 10⁴

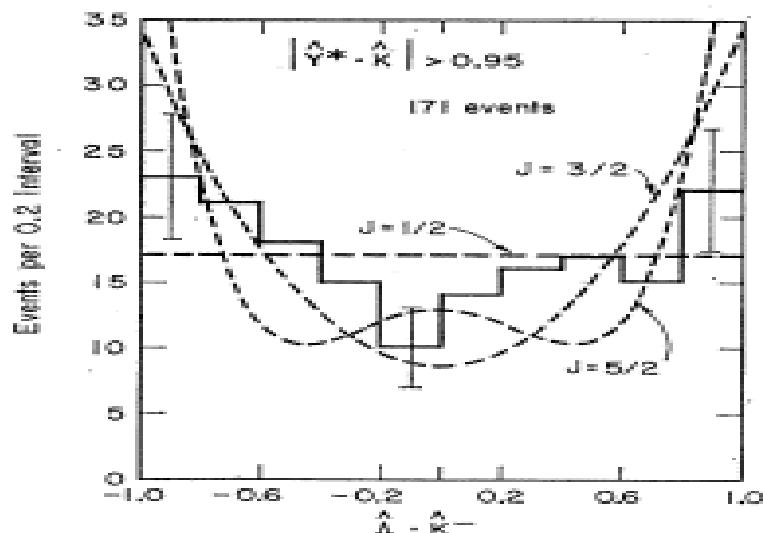
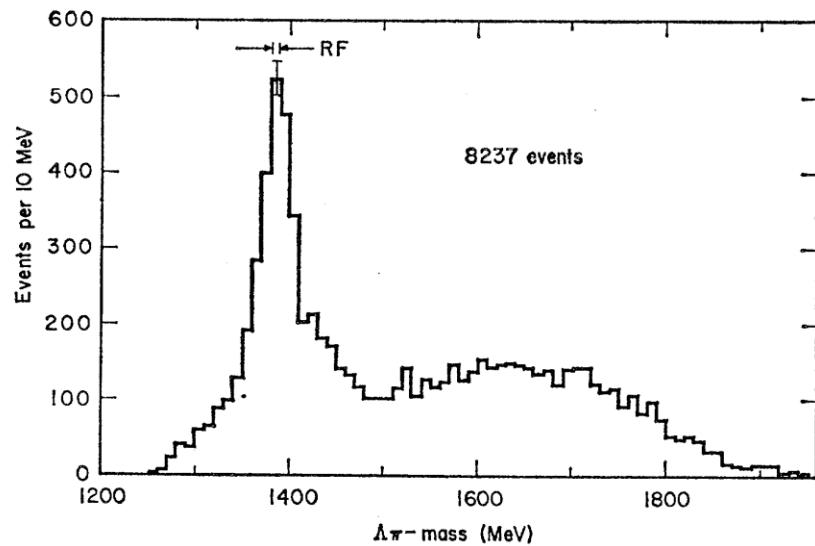
$\bar{p} \Delta(1232)^+$	3/2+	< 1	}	SU(3) breaking
$\bar{\Sigma}^- \Sigma(1385)^+$		3.1 ± 0.5		
$\bar{\Xi}^+ \Xi(1520)^-$		5.9 ± 1.5		

$\bar{p} N^*(1535)^+$	1/2-	10 ± 3	}	SU(3) allowed
$\bar{\Sigma}^- \Sigma(1360)^+$?		
$\bar{\Xi}^+ \Xi(1530)^-$?		

It is very important to check whether under the $\Sigma(1385)$ and $\Xi(1520)$ peaks there are 1/2- components ?

Evidence for the predicted $\Sigma^*(1/2^-)$

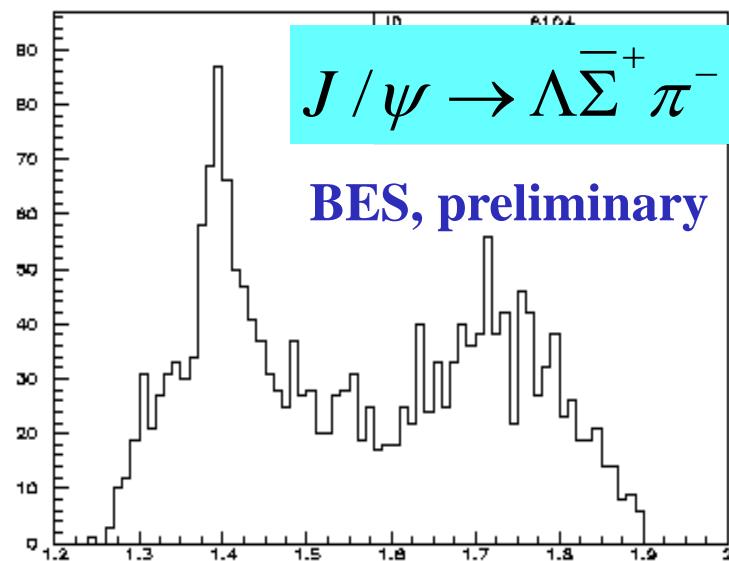
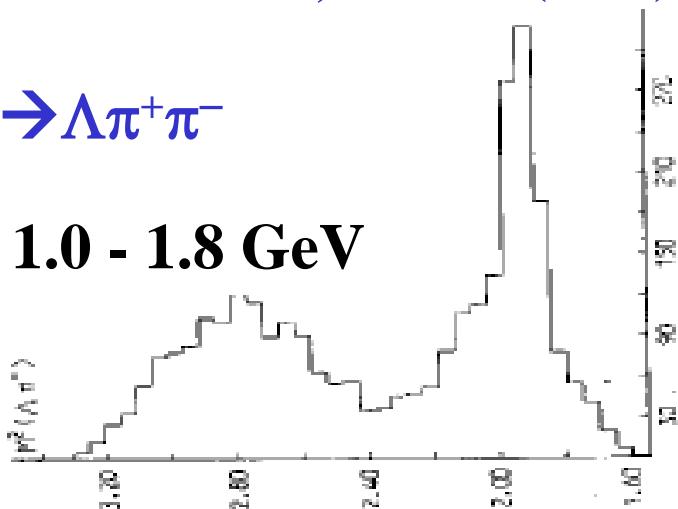
Huwe, PR181(1969)1824



Cameron et al., NPB143(1978)189

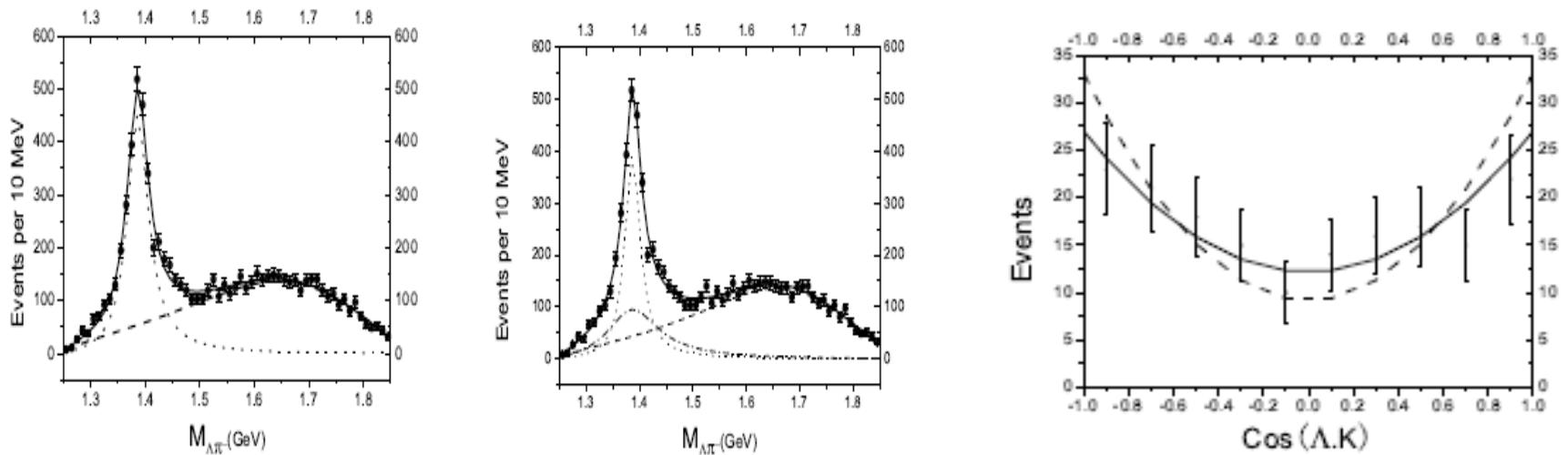


$$P_K = 1.0 - 1.8 \text{ GeV}$$



BES, NSTAR04

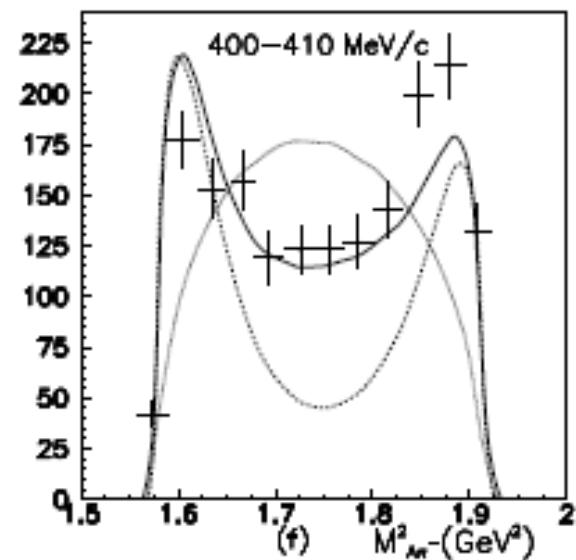
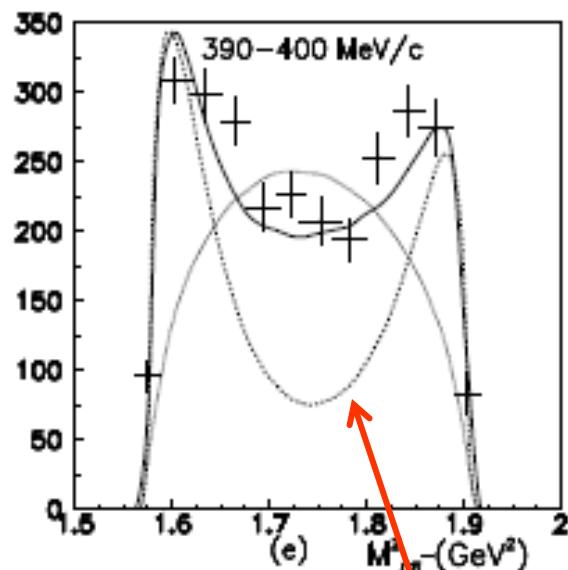
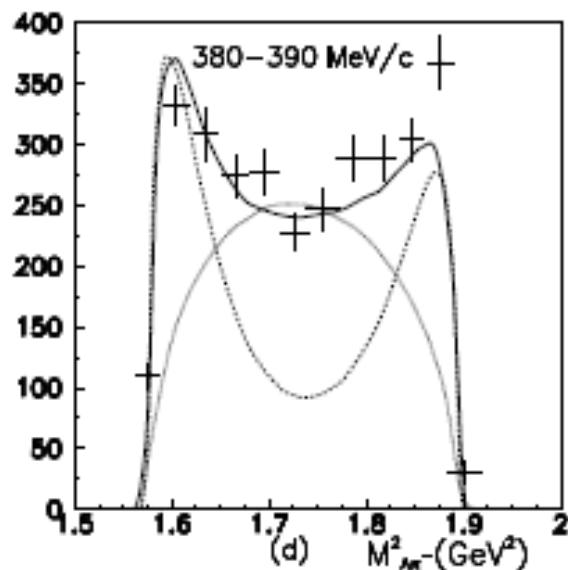
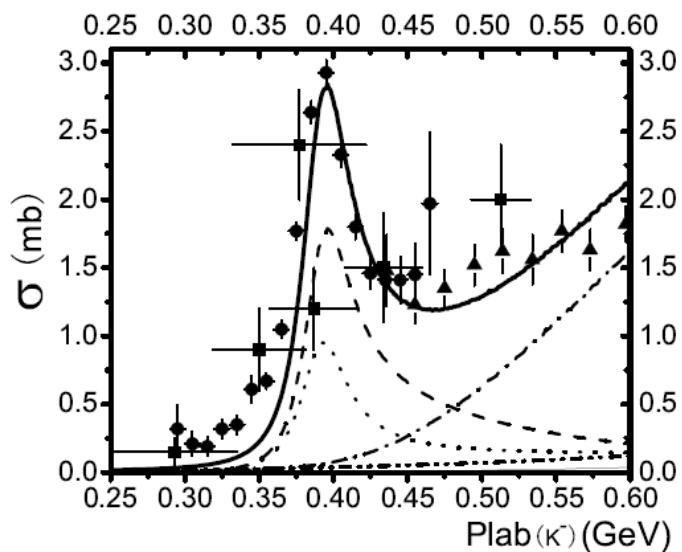
$M_{\pi\Lambda}$



$M_{\Sigma^*(3/2)}$	$\Gamma_{\Sigma^*(3/2)}$	$M_{\Sigma^*(1/2)}$	$\Gamma_{\Sigma^*(1/2)}$	χ^2/ndf (Fig.1)	χ^2/ndf (Fig.2)
Fit1 1385.3 ± 0.7	46.9 ± 2.5			68.5/54	10.1/9
Fit2 $1386.1^{+1.1}_{-0.9}$	$34.9^{+5.1}_{-4.9}$	$1381.3^{+4.9}_{-8.3}$	$118.6^{+55.2}_{-35.1}$	58.0/51	3.2/9

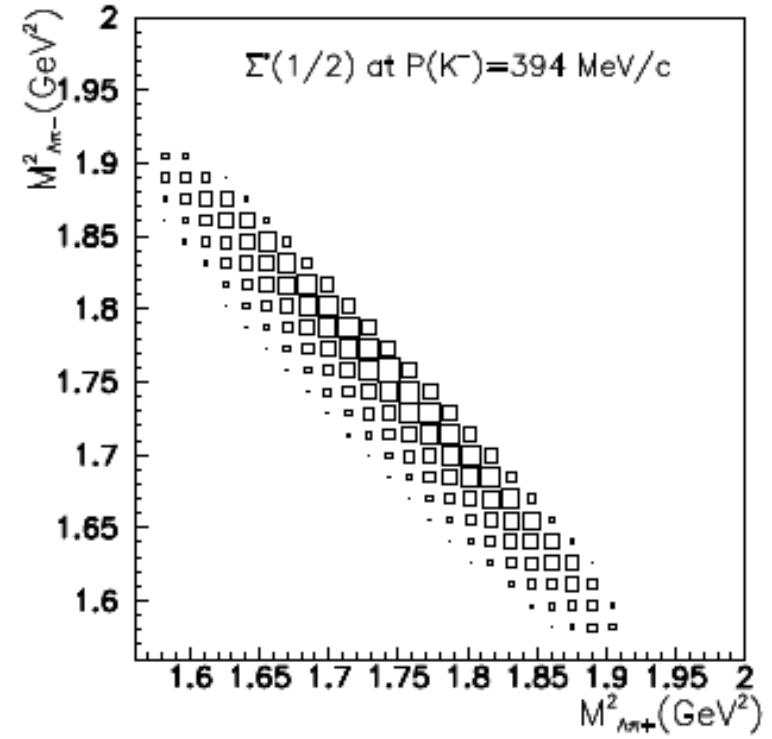
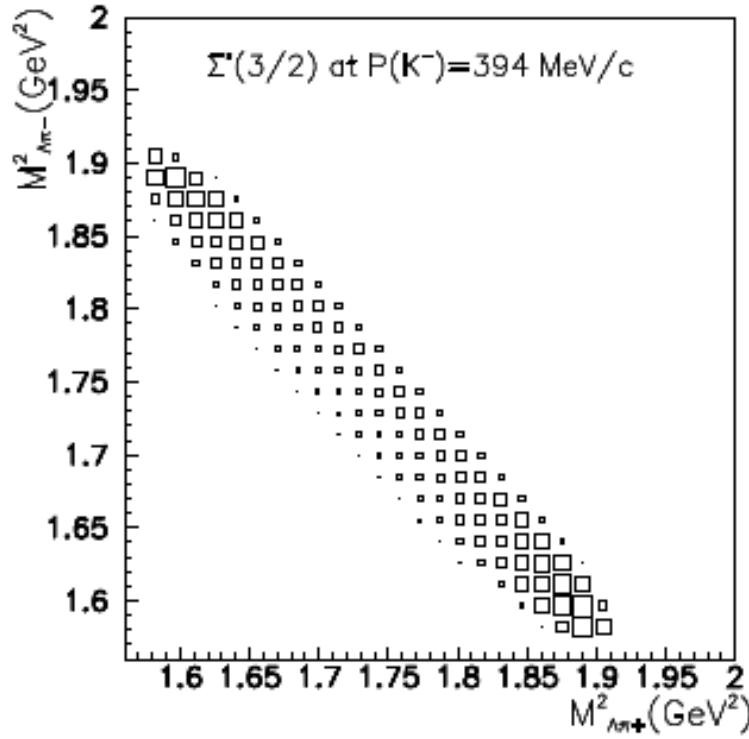
$$K^- p \rightarrow \Lambda^* \rightarrow \Sigma_{3/2}^{*-} \pi^+ \rightarrow \Lambda \pi^+ \pi^-$$

$$K^- p \rightarrow \Lambda^* \rightarrow \Sigma_{1/2}^{*-} \pi^+ \rightarrow \Lambda \pi^+ \pi^-$$

$$P_K \approx 0.4 \text{ GeV}$$


$\Sigma^*(3/2^+) \text{ only}$

J.J.Wu, S.Dulat, B.S.Zou, Phys. Rev. C81 (2010) 045210



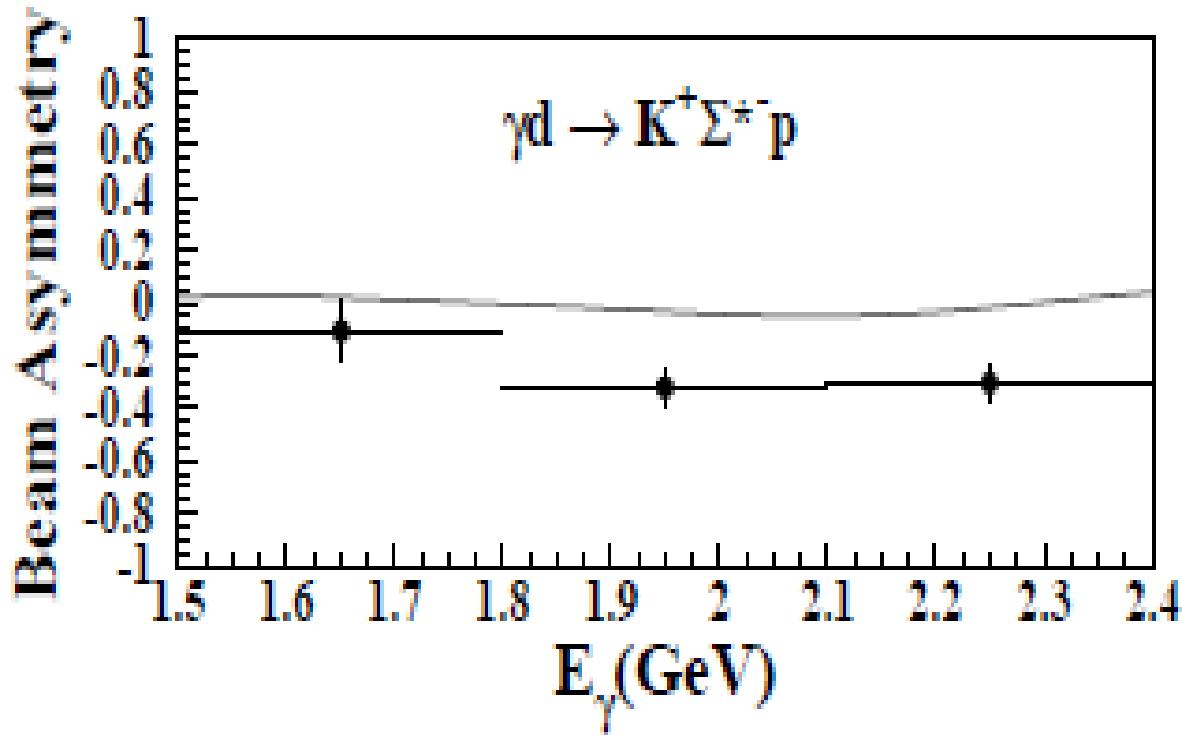
$\Sigma^*(3/2^+) & \Sigma^*(1/2^-) \rightarrow$ different Dalitz plots & mass spectra

Both are needed to reproduce the data !

Other evidence: failed to reproduce data with $\Sigma^*(1385)$

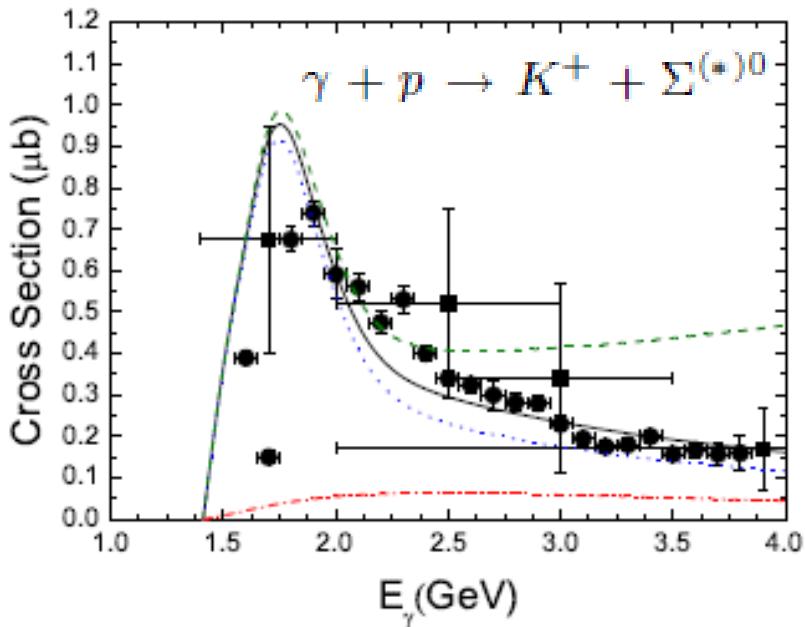
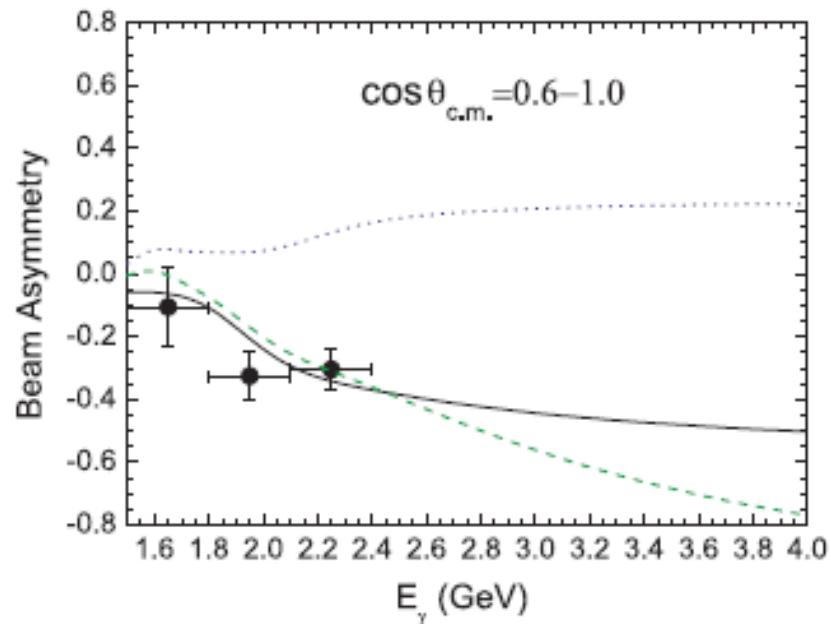
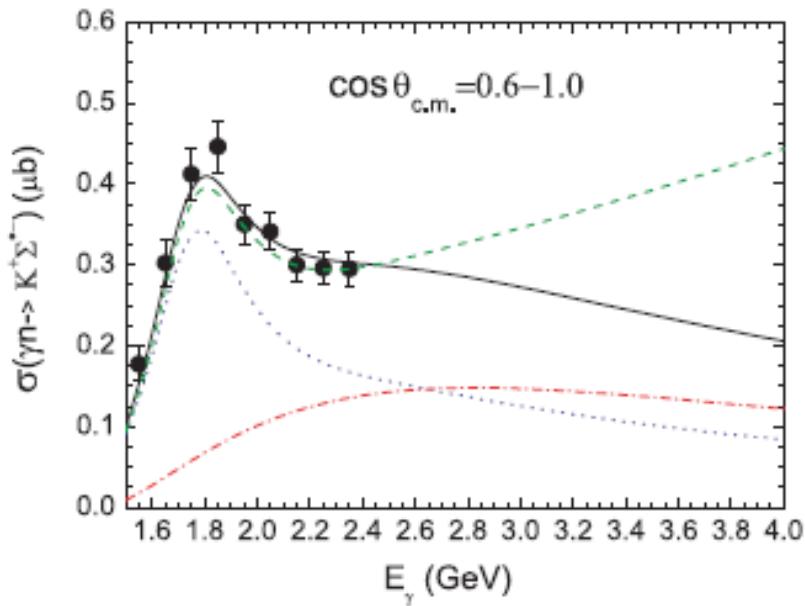
LEPS, PRL102(2009)012501

Y. Oh, C. M. Ko, and K. Nakayama, PRC77(2008) 045204



Something new ? $\Sigma^*(1/2^-)$?

P.Gao, J.J.Wu, B.S.Zou, Phys. Rev. C 81 (2010) 055203



dot lines: $\Sigma^*(3/2^+)$ with $h=1.00$
dashed : $\Sigma^*(3/2^+)$ with $h=1.11$
solid: including $\Sigma^*(1/2^-)$

P.Gao, J.J.Wu, B.S.Zou,
Phys. Rev. C 81 (2010) 055203

5. Summary and Prospects

$J/\psi \rightarrow \bar{N} \pi N$: $N^*(1440)$ peak at 1360 MeV
 $N^*(2065) \ 1/2+ \ \& \ 3/2+$

$J/\psi \rightarrow \bar{N} \eta N$: $N^*(1535)$

$J/\psi \rightarrow \bar{\Lambda} KN$: N_x or $N^*(1535)$?
Large [ud][us] \bar{s} ?

High statistics and PWA are needed !

Prospects :

2001 58M J/ Ψ

2009 BEPCII double ring upgrade



Completing N^* , Λ^* , Σ^* , Ξ^* spectra, deducing ΨBB^* couplings which provide a new way for exploring baryon structure.

CEBAF, CLEO-c, JPARC(kaon beam), COSY, CSR, PANDA, ...
competition & complementary



? Important Discoveries

Thanks !